



Optical Fibers and Their Biomedical Applications

Govind P. Agrawal

The Institute of Optics
University of Rochester
Rochester, New York, USA

©2013 G. P. Agrawal



Back

Close

Introduction

- Optical Fibers have found a variety of applications.
- High-capacity fiber-Optic links installed around the world make the magic of Internet and Telecommunications possible.
- Biomedical applications of fibers are increasing with the advent of photonic crystal and other microstructured fibers.
- Fibers act as optical wires and transport light just as electrical wires transport electricity.
- Nonlinear effects inside fibers provide a variety of new applications.
- This seminar provides a general overview of optical fibers with emphasis on their biomedical applications.
- Physical mechanism behind optical fibers: Total Internal Reflection (TIR).



2/40



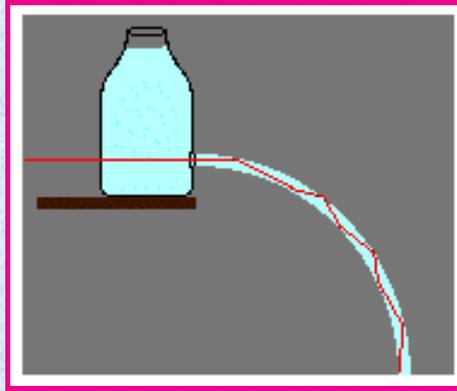
Back

Close

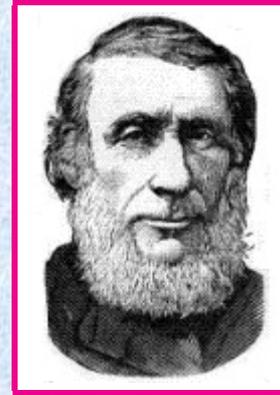
Historical Evolution



Daniel Colladon



Experimental Setup



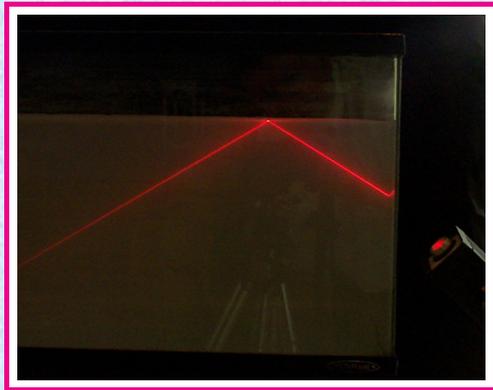
John Tyndall

- TIR is attributed to John Tyndall (1854 experiment in London).
- The book *City of Light* by Jeff Hecht (1999) traces history of TIR.
- First demonstration in Geneva in 1841 by Daniel Colladon (Comptes Rendus, vol. 15, pp. 800-802, Oct. 24, 1842).
- Light remained confined to a falling stream of water.



Historical Details

- Tyndall repeated the experiment in a 1854 lecture at the suggestion of Faraday (but Faraday could not recall the original name).
- Tyndall's name got attached to TIR because he described the experiment in his popular book *Light and Electricity* (around 1860).
- Colladon published an article **The Colladon Fountain** in 1884 to claim credit but it didn't work (La Nature, Scientific American).



A fish tank and a laser pointer can be used to demonstrate the phenomenon of total internal reflection.



4/40



Back

Close

Early TIR Applications



- Luminous fountains used TIR during the “Exposition Universel” in Paris in November 1889.
- Photograph from La Nature, published in 1889 (Fig. 3 on p. 593)



5/40

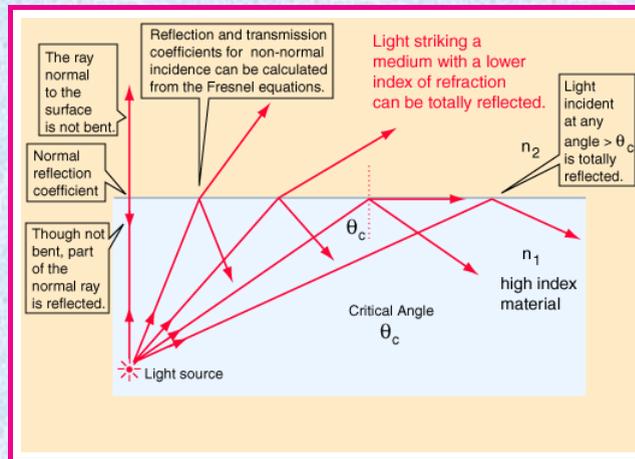


Back

Close

Total Internal Reflection

- Refraction of light at a dielectric interface is governed by Snell's law: $n_1 \sin \theta_i = n_2 \sin \theta_t$ (around 1620).
- When $n_1 > n_2$, light bends away from the normal ($\theta_t > \theta_i$).
- At a critical angle $\theta_i = \theta_c$, θ_t becomes 90° (parallel to interface).
- Total internal reflection occurs for $\theta_i > \theta_c$.



Cladded Glass Fibers

- Unclad glass fibers (diameter < 1 mm) were made during 1920s.
- A fiber bundle for transmitting images patented by C. W. Hansell and used by H. Lamm of Germany in 1930.
- Bare glass fibers suffered from several practical issues.
- Van Heel in Holland coated glass fibers with silver but no light emerged from them (because of losses during multiple reflections).
- Brian O'Brien (Univ. of Rochester) suggested to van Heel the use of dielectric cladding in October 1951.
- By April 1952, Van Heel transmitted images over 50 cm using a bundle of 400 glass fibers with plastic cladding.



7/40



Back

Close

Early Fiber Applications

- Hirschowitz wanted to use optical fibers for making a flexible gastroscope in 1954.
- Glass cladding was used in 1956 by L. Curtiss at Univ. of Michigan.
- He used a glass rod inside a glass tube of lower refractive index.
- Such a fiber was used to build a gastroscope in Feb. 1957.
- First commercial fiber gastroscope became available in 1960 and soon after it was adopted by hospitals worldwide.
- Other early applications included cystoscopes, image processing, and toys. Losses limited fiber lengths to a few meters.
- Narinder Kapany popularized fiber optics with his 1960 Scientific American article, followed with a book.



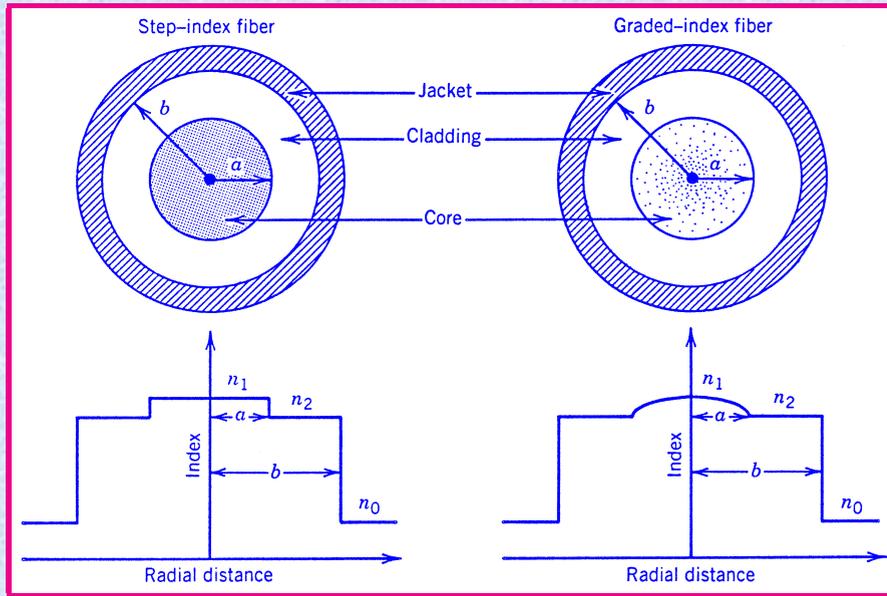
8/40



Back

Close

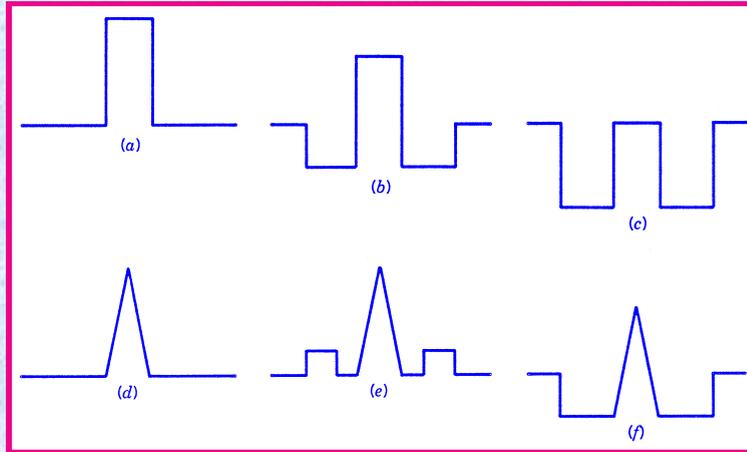
Modern Optical Fibers



- Contain a central core surrounded by a lower-index cladding
- Two-dimensional waveguides with cylindrical symmetry
- Graded-index fibers: Refractive index varies inside the core.



Fiber Design



- Core doped with GeO_2 ; cladding with fluorine.
- Index profile rectangular for standard fibers.
- Triangular index profile for dispersion-shifted fibers.
- Raised or depressed cladding for dispersion control.



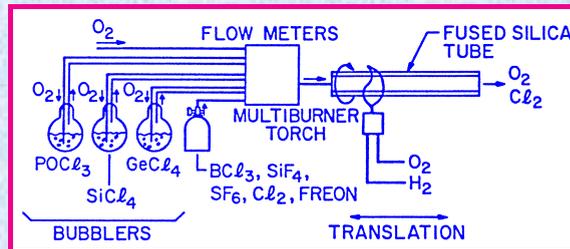
Silica Fibers

Two-Stage Fabrication

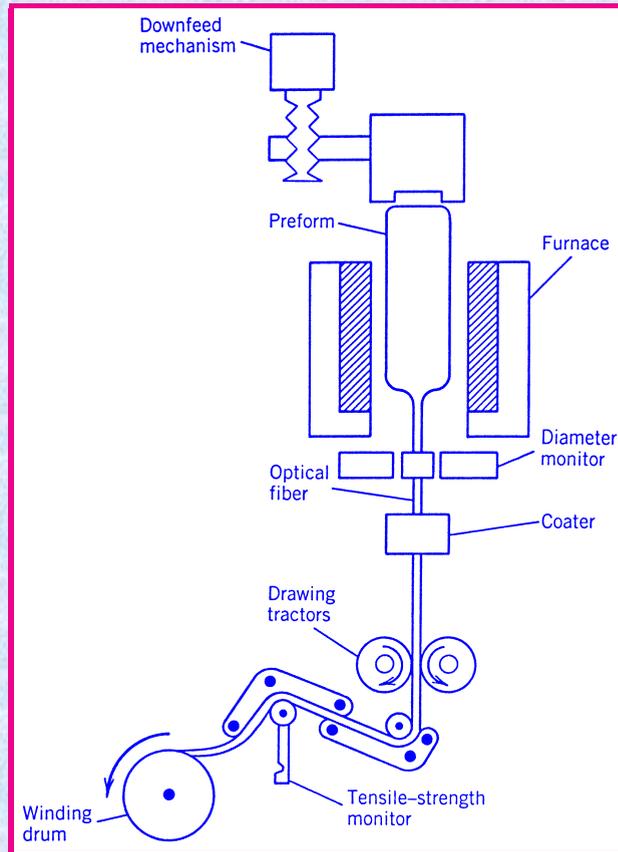
- **Preform:** Length 1 m, diameter 2 cm; correct index profile.
- Preform is drawn into fiber using a draw tower.

Preform Fabrication Techniques

- Modified chemical vapor deposition (MCVD).
- Outside vapor deposition (OVD).
- Vapor Axial deposition (VAD).



Fiber Draw Tower



Fiber Properties

- Light guided over long lengths through total internal reflection.
- Propagation within the fiber depends on the number of modes supported by that fiber.
- One must solve Maxwell's equations to find these spatial modes.
- Important fiber parameter: $V = (2\pi/\lambda)a\sqrt{n_1^2 - n_2^2}$.
- Each mode has an effective refractive index: $n_2 < \bar{n} < n_1$.
- Only a single guided mode exists in fibers designed with $V < 2.4$.
- Typically $a = 5 \mu\text{m}$ for $\lambda > 1.3 \mu\text{m}$ for single-mode fibers.
- Nonlinear effects are enhanced in narrow-core fibers because of a reduced effective mode area.



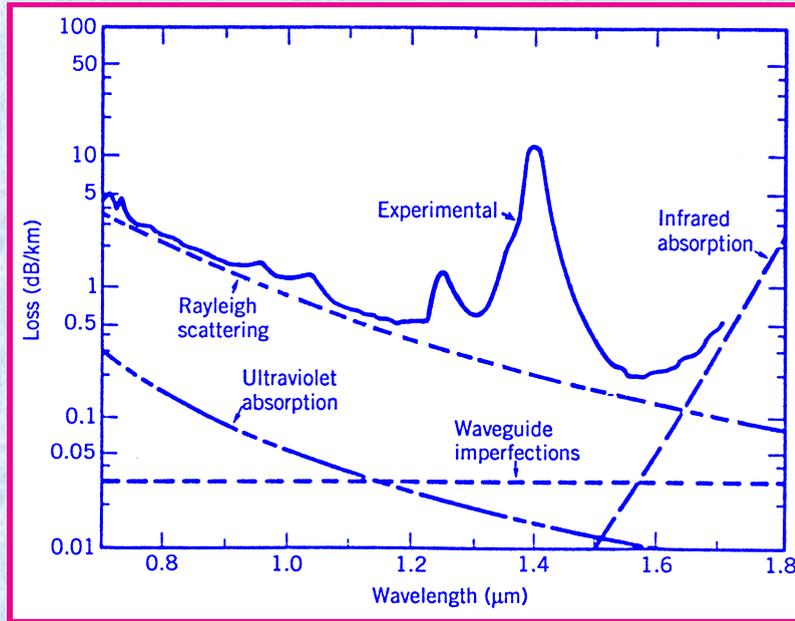
13/40



Back

Close

Fiber Losses



Definition: $P_{\text{out}} = P_{\text{in}} \exp(-\alpha L)$, α (dB/km) = 4.343 α .

- Water peak near 1400 nm reduced considerably in modern fibers.
- Rayleigh scattering is the dominant loss mechanism (varies as λ^{-4})



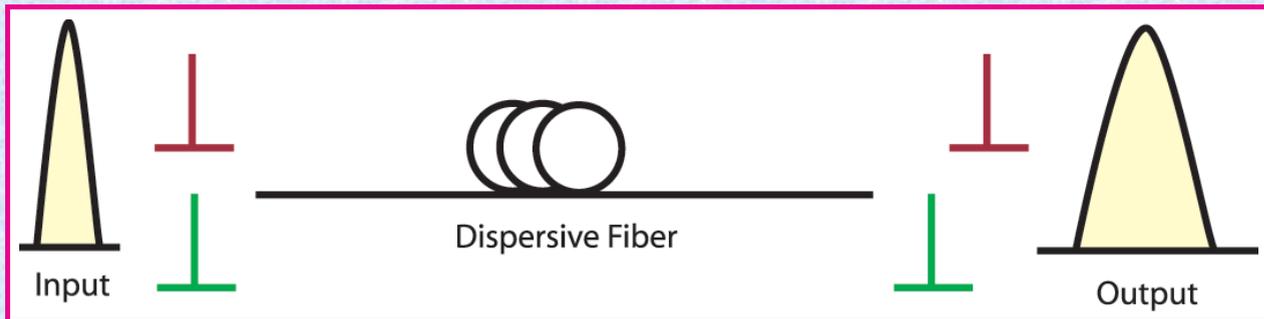
Fiber Dispersion

- Origin: Frequency dependence of the mode index $\bar{n}(\omega)$:

$$\beta(\omega) = \bar{n}(\omega)\omega/c = \beta_0 + \beta_1(\omega - \omega_0) + \frac{1}{2}\beta_2(\omega - \omega_0)^2 + \dots,$$

where ω_0 is the carrier frequency of optical pulse.

- *Transit time* for a fiber of length L : $T = L/v_g = L(d\beta/d\omega)$.
- Different frequency components travel at different speeds and arrive at different times at output end (pulse broadening).



Fiber Dispersion

- Pulse broadening governed by group-velocity dispersion:

$$\Delta T = \frac{dT}{d\omega} \Delta\omega = \frac{d}{d\omega} \left(\frac{L}{v_g} \right) = L\beta_2 \Delta\omega,$$

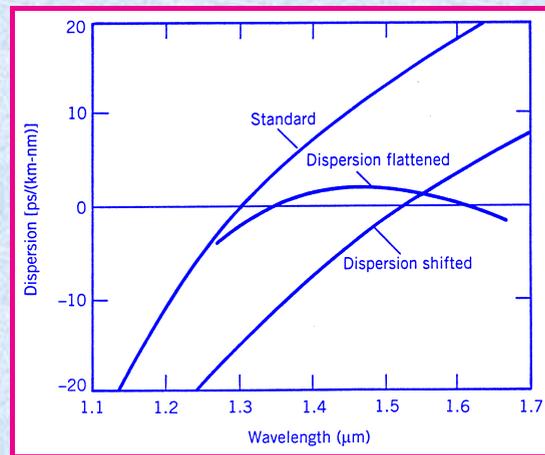
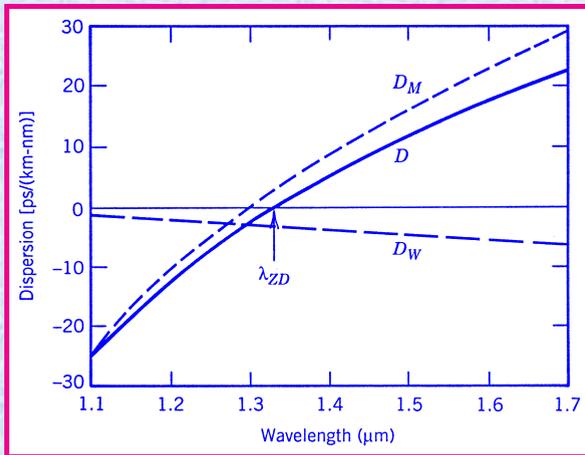
$\Delta\omega$ is pulse bandwidth and L is fiber length.

- Second-order dispersion: $\beta_2 = \frac{d}{d\omega} \left(\frac{1}{v_g} \right) = \frac{d^2\beta}{d\omega^2}$ (ps²/km)
- Alternate definition: $D = \frac{d}{d\lambda} \left(\frac{1}{v_g} \right) = -\frac{2\pi c}{\lambda^2} \beta_2$ (ps/km/nm)
- Third-order Dispersion, governed by β_3 , becomes important for femtosecond pulses. It is related to dispersion slope $\frac{dD}{d\lambda}$.



Waveguide Dispersion

- Mode index $\bar{n}(\omega) = n_1(\omega) - \delta n_W(\omega)$.
- Material dispersion D_M results from $n_1(\omega)$ (index of silica).
- Waveguide dispersion D_W results from $\delta n_W(\omega)$ and depends on core size and dopant distribution.
- Total dispersion $D = D_M + D_W$ can be controlled.



Major Nonlinear Effects

- Self-Phase Modulation (SPM)
- Cross-Phase Modulation (XPM)
- Four-Wave Mixing (FWM)
- Stimulated Brillouin Scattering (SBS)
- Stimulated Raman Scattering (SRS)

Origin of Nonlinear Effects in Optical Fibers

- Third-order nonlinear susceptibility $\chi^{(3)}$.
- Real part leads to SPM, XPM, and FWM.
- Imaginary part leads to SBS and SRS.



18/40



Back

Close

Self-Phase Modulation (SPM)

- Refractive index depends on intensity as

$$n(I) = n + n_2 I(t).$$

- $n_2 = 2.6 \times 10^{-20} \text{ m}^2/\text{W}$ for silica fibers.
- Propagation constant: $\beta' = \beta + k_0 \bar{n}_2 (P/A_{\text{eff}}) \equiv \beta + \gamma P$.
- Nonlinear parameter: $\gamma = 2\pi \bar{n}_2 / (A_{\text{eff}} \lambda)$.
- Nonlinear Phase shift:

$$\phi_{\text{NL}} = \int_0^L \gamma P_0 e^{-\alpha z} dz = \gamma P_0 L_{\text{eff}}, \quad L_{\text{eff}} = (1 - e^{-\alpha L}) / \alpha.$$

- Optical field modifies its own phase (SPM).
- SPM leads to spectral broadening of optical pulses.



19/40



Back

Close

Nonlinear Schrödinger Equation

- Dispersive and Nonlinear effects are studied by solving a simple equation derived from Maxwell's equations.
- This equation is known as the Nonlinear Schrödinger Equation:

$$\frac{\partial A}{\partial z} + \frac{i\beta_2}{2} \frac{\partial^2 A}{\partial t^2} = i\gamma|A|^2A - \frac{\alpha}{2}A.$$

- Nonlinear parameter: $\gamma = 2\pi\bar{n}_2/(A_{\text{eff}}\lambda)$.
- Fibers with large A_{eff} are useful for telecommunications.
- Highly nonlinear fibers are made with reduced A_{eff} .
- Nonlinear effects are enhanced considerably in photonic crystal and other microstructured fibers.



20/40



Back

Close

Cross-Phase Modulation (XPM)

- Refractive index seen by one wave depends on the intensity of other co-propagating waves.

- Nonlinear phase shifts for two waves take the form

$$\phi_a^{\text{NL}} = \gamma_a L_{\text{eff}}(P_a + 2P_b), \quad \phi_b^{\text{NL}} = \gamma_b L_{\text{eff}}(P_b + 2P_a).$$

- The second term with a factor of 2 is due to XPM.
- Mathematically, XPM effects are governed by two coupled NLS equations.
- XPM is beneficial for applications such as optical switching, wavelength conversion, etc.
- XPM is a limiting factor for modern telecommunication systems.



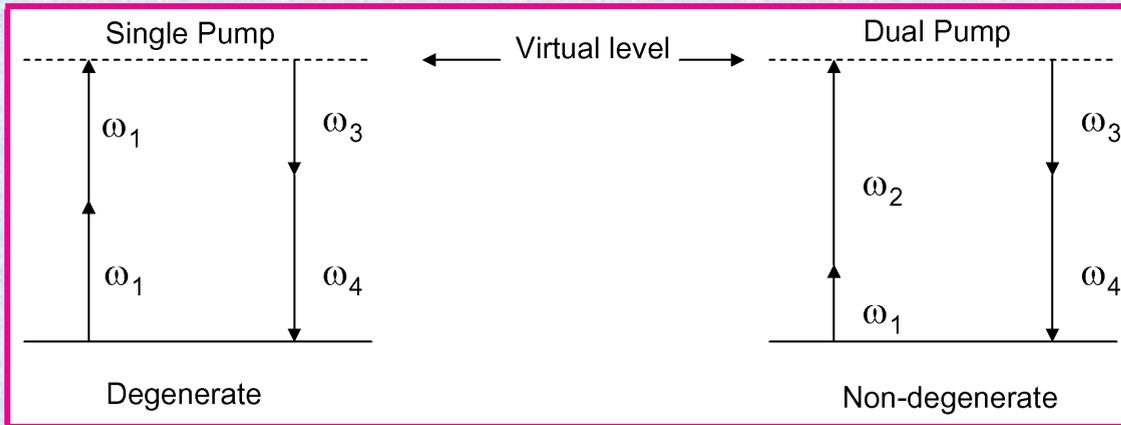
21/40



Back

Close

Four-Wave Mixing

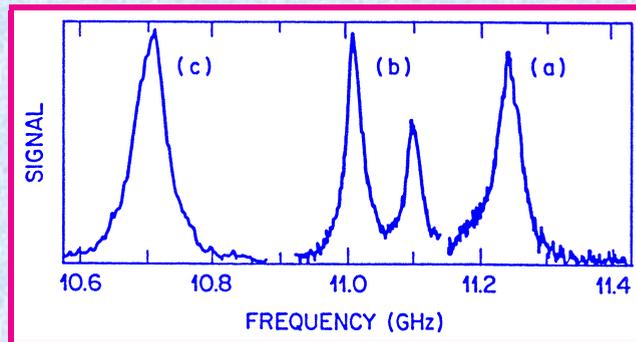


- FWM converts two photons from one or two pump beams into two new frequency-shifted photons.
- Energy conservation: $\hbar\omega_1 + \hbar\omega_2 = \hbar\omega_3 + \hbar\omega_4$.
- Momentum conservation leads to phase-matching condition.
- Useful for parametric amplification and wavelength conversion.



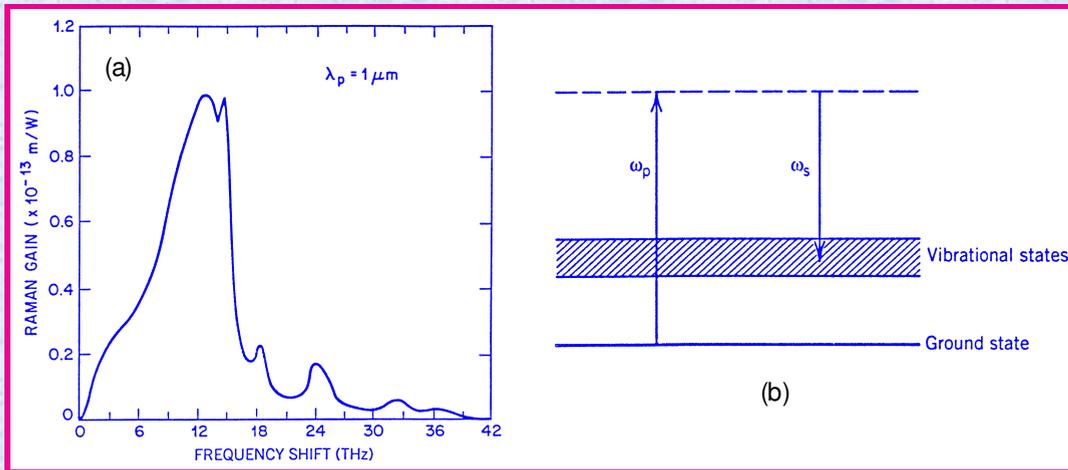
Brillouin Scattering

- Scattering of light from acoustic waves (electrostriction).
- Energy and momentum conservation laws require $\Omega_A = \omega_p - \omega_s$ and $\mathbf{k}_A = \mathbf{k}_p - \mathbf{k}_s$.
- Becomes a stimulated process at high input power levels.
- Most of power reflected once threshold is reached (<10 mW for fiber lengths >10 km).
- Brillouin gain has a narrow Lorentzian spectrum ($\Delta\nu \sim 20$ MHz).

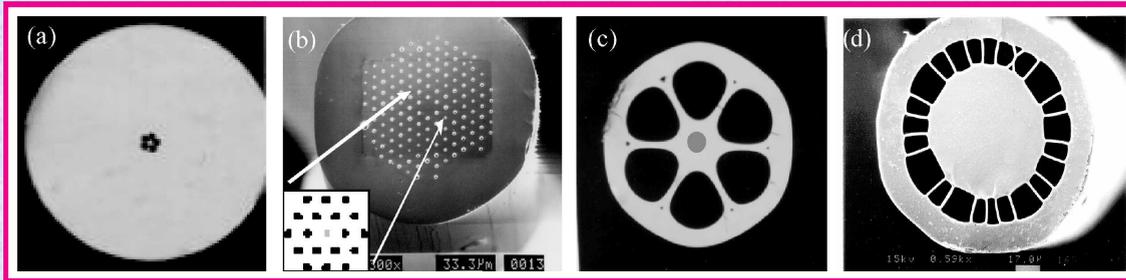


Stimulated Raman Scattering

- Scattering of light from vibrating molecules.
- Scattered light shifted in frequency.
- Raman gain spectrum extends over 40 THz.
- Raman shift at Gain peak: $\Omega_R = \omega_p - \omega_s \sim 13$ THz).



Microstructured Fibers



(Eggleton et al, Opt. Exp. **9**, 698, 2001)

- A narrow core is surrounded by a silica cladding with air holes.
- Photonic crystal fibers have multiple rings of holes.
- Number of air holes varies from structure to structure.
- Hole size varies from 0.5 to 5 μm depending on the design.
- Nonlinear effects are enhanced considerably (highly nonlinear fibers).
- Useful for supercontinuum generation among other things.



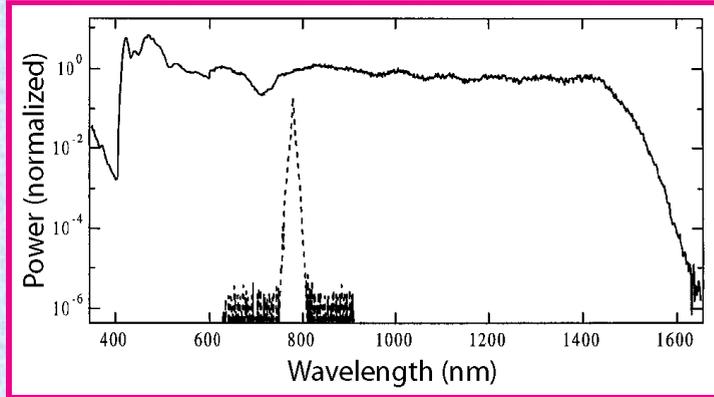
25/40



Back

Close

Supercontinuum Generation



(Ranka et al., Opt. Lett. **25**, 25, 2000)

- Output spectrum generated in a 75-cm section of microstructured fiber using 100-fs pulses with 0.8 pJ energy.
- Supercontinuum extends from 400 to 1600 nm.
- It is also relatively flat over the entire bandwidth.
- Useful in biomedical imaging as a broadband source.



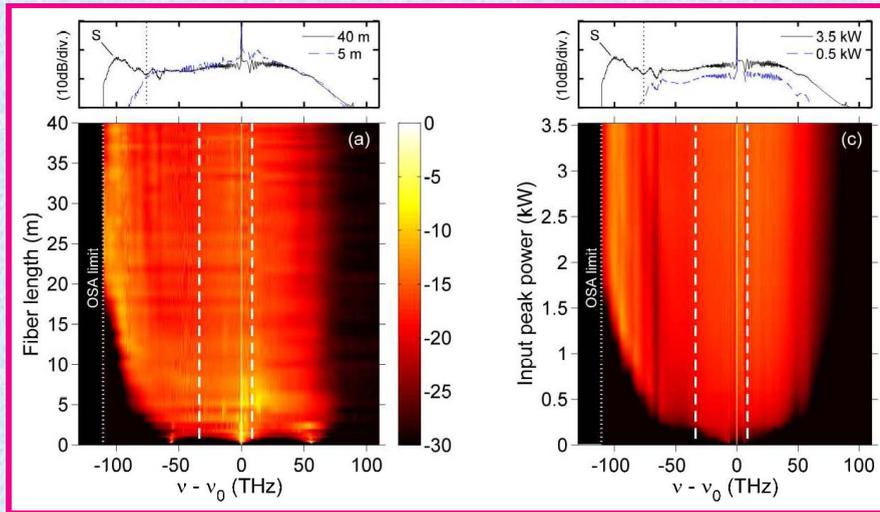
26/40



Back

Close

Dispersion-Tailored Fibers



(Barviau et al., Opt. Exp. **17**, 7392, 2009)

- SC spectra as a function of fiber length and input peak power.
- A PCF with two zero-dispersion wavelengths was employed.
- Supercontinuum is relatively flat over the entire bandwidth.



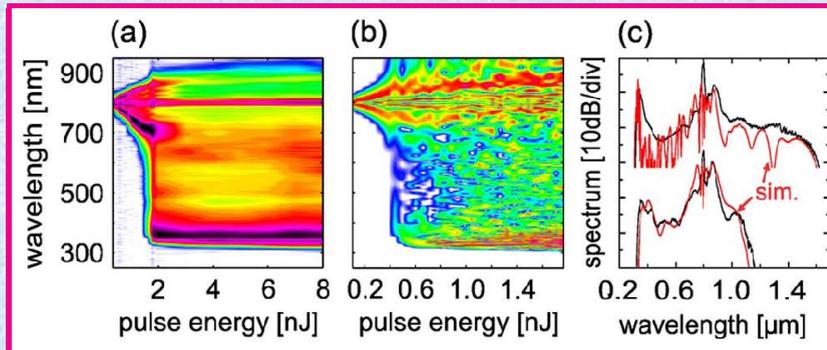
27/40



Back

Close

Tapered Photonic Crystal Fibers



(Stark et al., Opt. Lett. **37**, 770, 2012)

- Experimental (a) and simulated (b) SC spectra when 110-fs pulses launched into a tapered PCF.
- (c) SC spectra at input pulse energies of 2 and 5 nJ.
- Core diameter tapered form $4.5 \mu\text{m}$ to $0.6 \mu\text{m}$ over 1 cm.
- Tapering helps to extend the supercontinuum into the UV region.



28/40



Back

Close

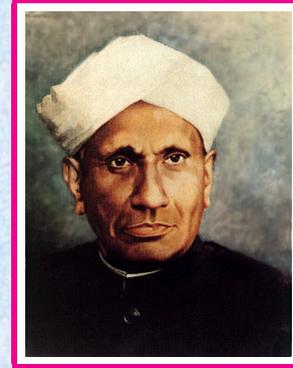
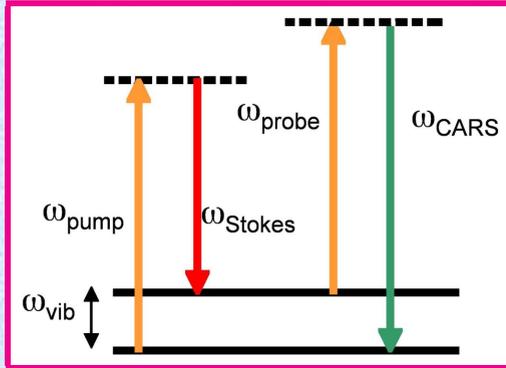
Supercontinuum-Based Biomedical Imaging



- Several companies sell fiber-based supercontinuum sources (NKT Photonics, Fianium, Koheras, Leukos, etc.).
- This has led to their use in biomedical imaging.
- Imaging techniques are known by a variety of names.
- I focus on 3 techniques: CARS microscopy; STED microscopy; optical coherence tomography.



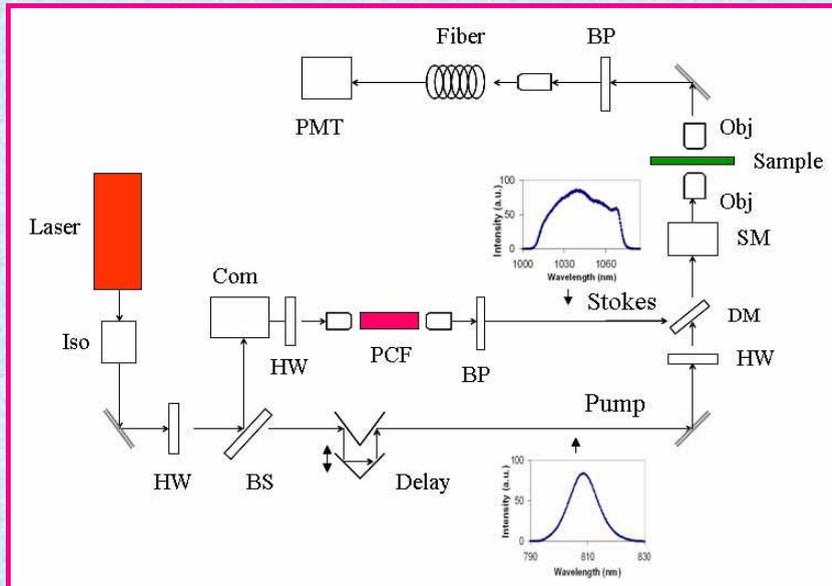
CARS Nonlinear Process



- Coherent anti-Stokes Raman scattering (CARS) is a well-known nonlinear process (Maker and Terhune, Phys. Rev., 1965).
- Pump and Stokes beam at ω_p and ω_s drive coherently a vibrational resonance at the frequency $\omega_{\text{vib}} = \omega_p - \omega_s$ (optical phonons).
- CARS signal generated at $\omega_{\text{CARS}} = 2\omega_p - \omega_s$.
- CARS is a kind of Raman-enhanced four-wave mixing process.



CARS Microscopy



(Murugkar et al., Opt. Exp. **15**, 14028, 2007)

Laser pulses (65-fs) split to produce pump and Stokes beams.
Bandpass filter after the PCF selects the Stokes bandwidth.
Different Stokes frequencies excite different molecules in sample.



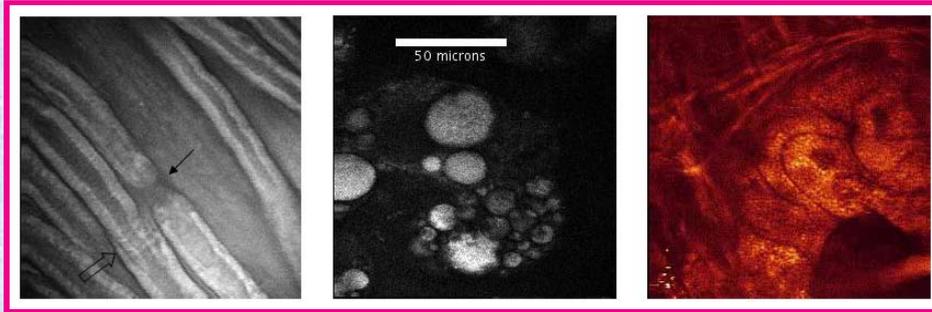
31/40



Back

Close

CARS Microscopy

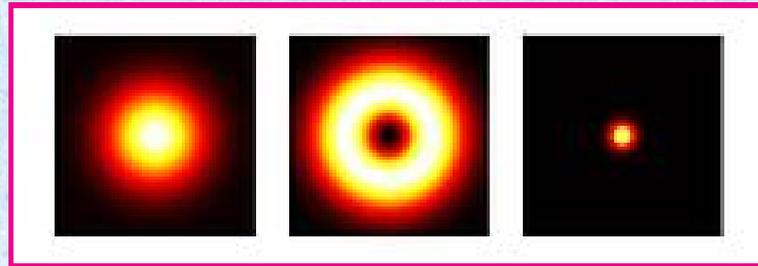
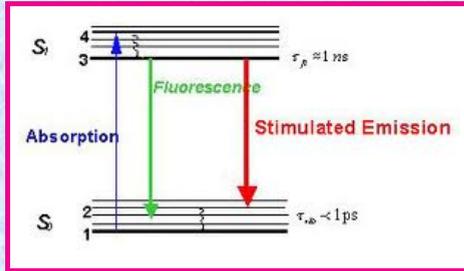


(Murugkar et al., Opt. Exp. **15**, 14028, 2007)

- Stokes pulses, broadened spectrally using a PCF, are sent to the sample together with pump pulses.
- Anti-Stokes signal generated inside the sample is used for microscopy.
- (a) Live rat dorsal root axon; (b) lipid droplets in a cell culture; (c) sebaceous gland in a mouse ear.
- Resolution is typically limited to 2-3 μm .



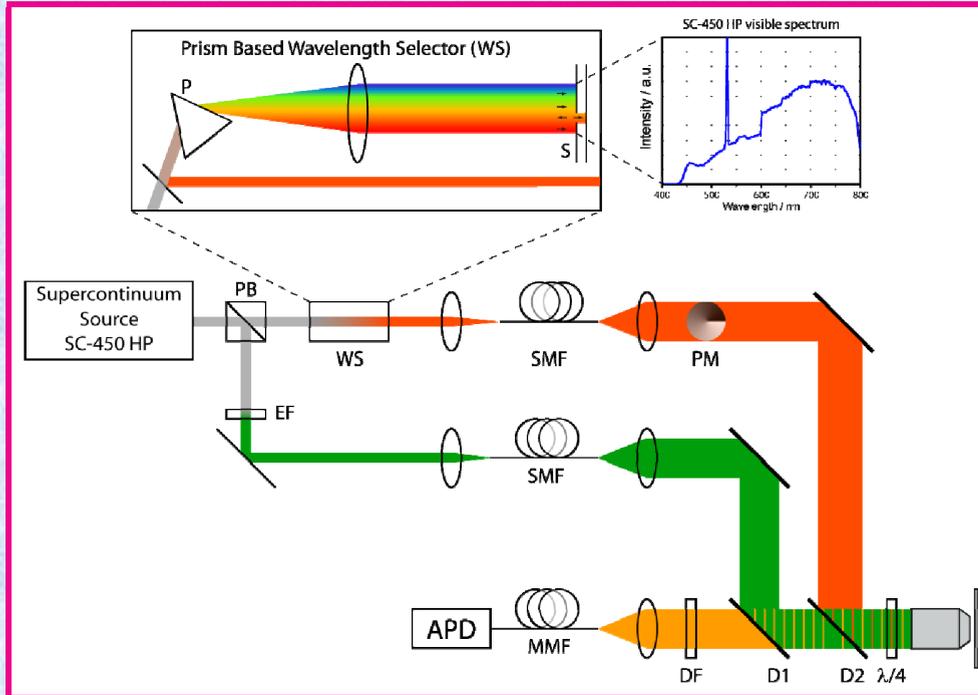
STED Microscopy



- Stimulated-emission depletion (STED) microscopy was first proposed in 1994 (Hell and Wichmann, Opt. Lett. **19**, 780, 1994).
- Fluorescence is suppressed in the off-center region using a second beam that removes excited molecules through stimulated emission.
- Nanoscale resolution ($\lambda/50$) realized by 2005 using a doughnut-shape STED pulsed beam.
- A fiber-based supercontinuum source was used by 2008.



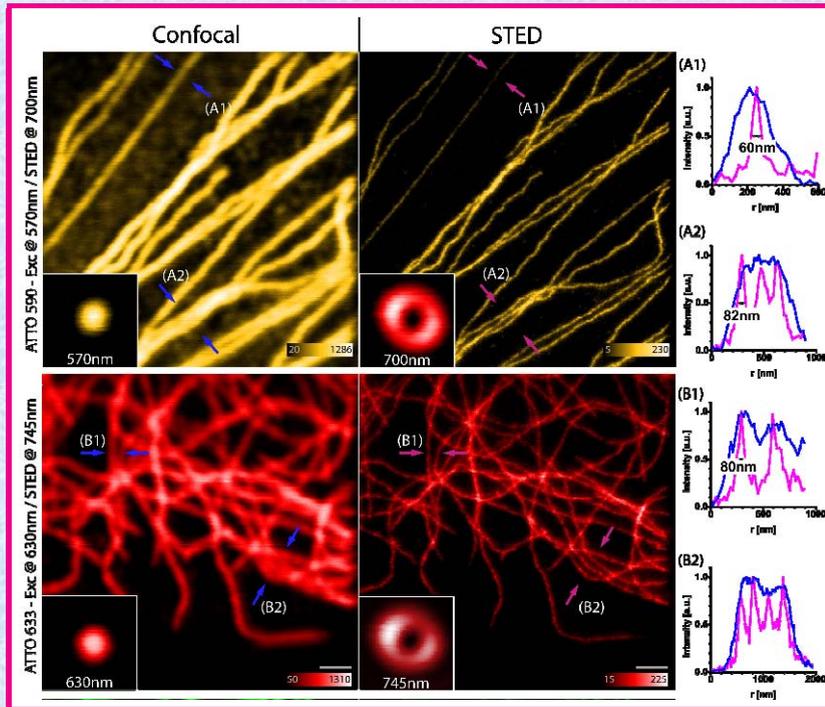
STED Setup



(Wildanger et al., Opt. Exp. **16**, 9614, 2008)



STED Results

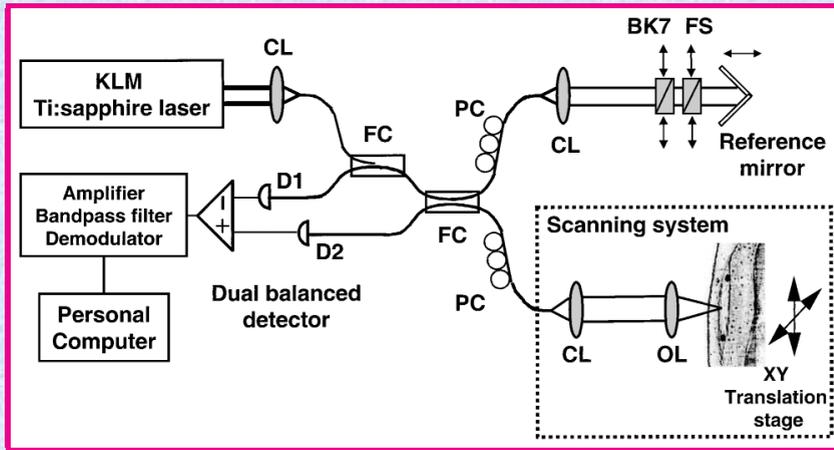


(Wildanger et al., Opt. Exp. **16**, 9614, 2008)

Immunolabeled tubulin fibers imaged at 570 nm and 630 nm.



Optical Coherence Tomography (OCT)



(Drexler et al., Opt. Lett. **24**, 1221, 1999)

- A linear imaging technique based on Michelson interferometry.
- Image resolution ($\Delta z = c\tau_c$) depends on the coherence time τ_c .
- Supercontinuum sources provide a resolution of $< 1 \mu\text{m}$.



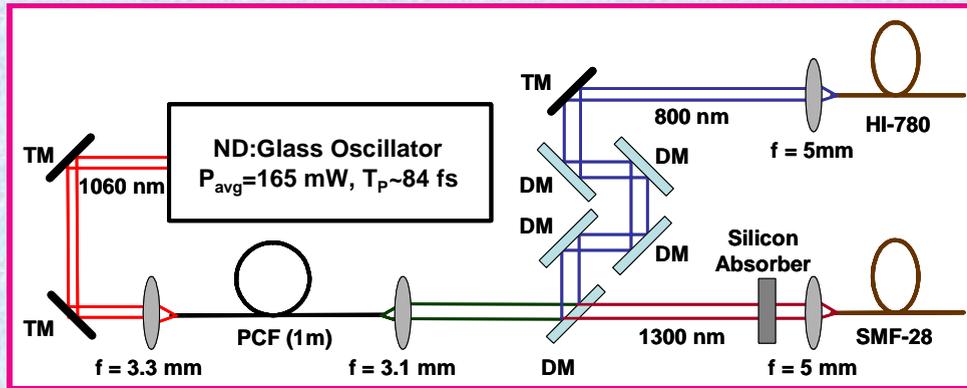
36/40



Back

Close

Dual-Band OCT



(Fujimoto et al., Opt. Exp. **14**, 1145, 2006)

- OCT is performed simultaneously using two spectral bands located near 800 and 1300 nm.
- Image resolution $< 3 \mu\text{m}$ at 800 and $< 5 \mu\text{m}$ at 1300 nm.
- Combined *in vivo* images of good quality possible.



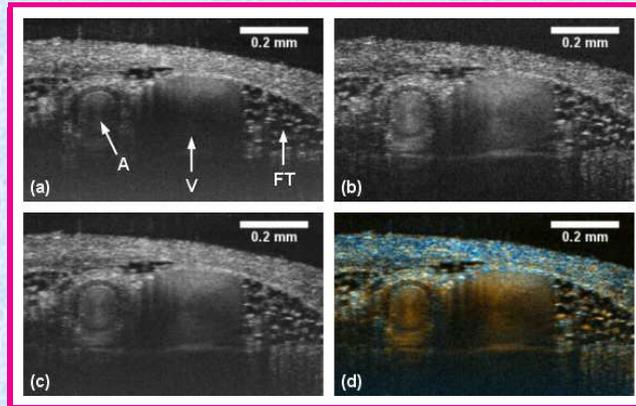
37/40



Back

Close

Dual-Band OCT

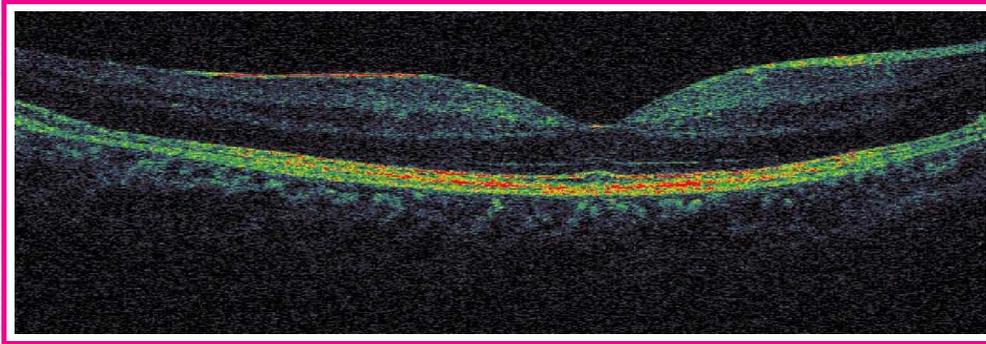


(Cimalla et al., Opt. Exp. **17**, 19486, 2009)

- Simultaneous *in vivo* scans of murine saphenous artery (A), vein (V) and perivascular fat tissue (FT) during the diastole.
- (a) Image at 800 nm; (b) same image at 1250 nm.
- (c) Compounded image of (a) and (b).
- (d) Color-encoded differential image.



High-Resolution OCT



(Nishiura et al., Jap. J. Appl. Phys. **49**, 012701, 2010)

- OCT *in vivo* image of human retina around fovea.
- Observed axial resolution was $2.1 \mu\text{m}$ in tissue.
- A Gaussian-shape 150-nm-wide supercontinuum was employed for this image.
- OCT is an established medical imaging technique. It is often used to image anterior segment of the eye or the retina.



39/40



Back

Close

Concluding Remarks

- Optical fibers were developed during the 1950s and used for biomedical applications during the 1960s.
- They became relevant for telecommunications after 1970 with the development of low-loss fibers.
- By 2000, more than 60 million kilometers of fiber was installed worldwide (on land and in the oceans).
- Biomedical applications of optical fibers are attracting attention in recent years.
- Nonlinear effects in optical fibers make it possible to create a supercontinuum whose bandwidth exceeds 100 THz.
- Such sources are useful for tissue tomography and nonlinear microscopy (biomedical imaging).



40/40



Back

Close