



# **Linear and Nonlinear Interference**

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# Outline

- ❑ **Introduction**
- ❑ **Wave superposition**
- ❑ **Radiation-matter interaction**
  - ❑ **Interference**
  - ❑ **Nonlinear optics**
- ❑ **Conclusions**

# Introduction

Electromagnetic waves exist in two independent polarizations and carry three quantities which can be transferred to matter:

- Energy
- Momentum
- Angular momentum.

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} + \omega t + \phi)}$$

The algebra of operating with EM fields is well defined (Jackson, Born Wolf) but has some hidden tricks

# Introduction

Over the years I have work both in Interferometry and Nonlinear Optics.

These are traditionally independent and different areas based on its scope.

But at the end I have found them to be basically equivalent.

# Wave Superposition

$$E = E_0 e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

$$E_1 = E_{01} e^{i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} \quad E_2 = E_{02} e^{i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)}$$

$$E_1 + E_2 = E_{01} e^{i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} + E_{02} e^{i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)}$$

# Wave Superposition: Interference

Every textbook follows the same steps:

$$E_1 + E_2 = E_{01} e^{i(\vec{k}_1 \cdot \vec{r} - \omega_1 t)} + E_{02} e^{i(\vec{k}_2 \cdot \vec{r} - \omega_2 t)}$$

$$I = \frac{1}{2} c \varepsilon_0 |E_1 + E_2|^2$$

$$= \frac{1}{2} c \varepsilon_0 \left| E_{01} e^{i\vec{k}_1 \cdot \vec{r}} + E_{02} e^{i\vec{k}_2 \cdot \vec{r}} \right|^2$$

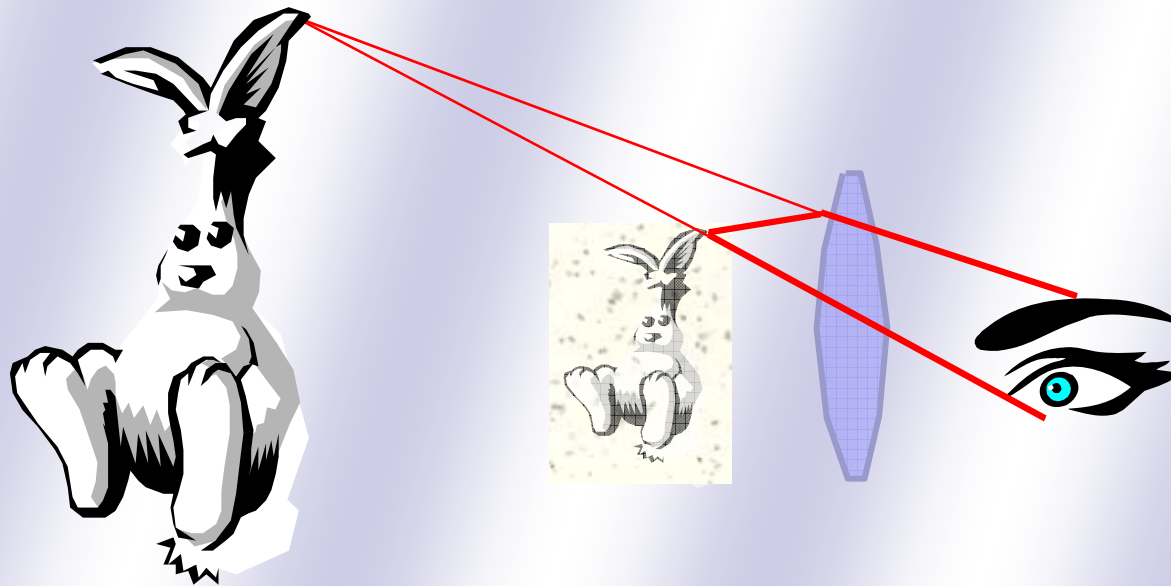
$$= \frac{1}{2} c \varepsilon_0 \left\{ |E_{01}|^2 + |E_{02}|^2 + |E_{01}| |E_{02}| \cos\left(\left(\vec{k}_1 - \vec{k}_2\right) \cdot \vec{r}\right) \right\}$$

**and so on ...**

# Interference to the limit

Even going one step further, image formation can be considered an interference phenomena.

The best “in focus image” is achieved when the rays coming from the source interfere constructively (somehow incorporated in physical optics math)

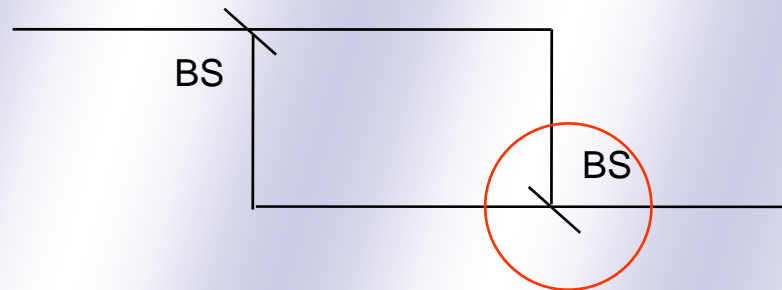


# Interference

The fundamental step is the observation of the interference pattern:

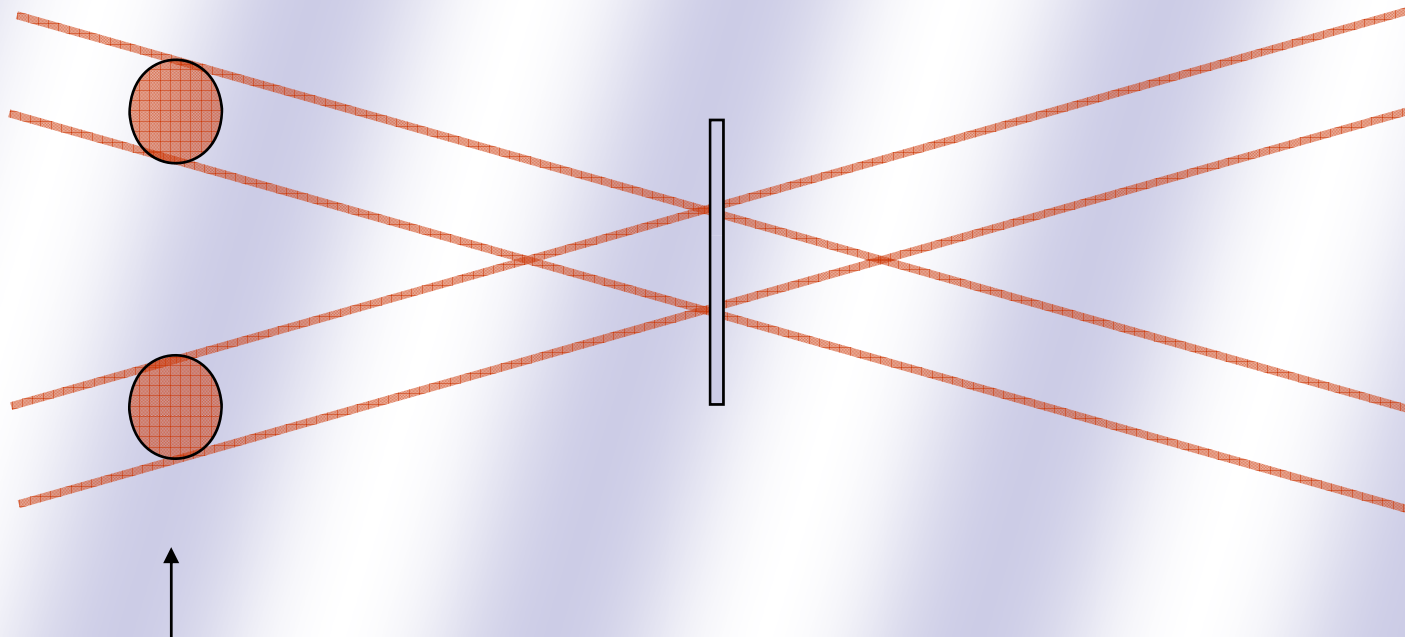
$$I = \frac{1}{2} c \varepsilon_0 |E_1 + E_2|^2$$

Lets see it in more detail on an interferometer:

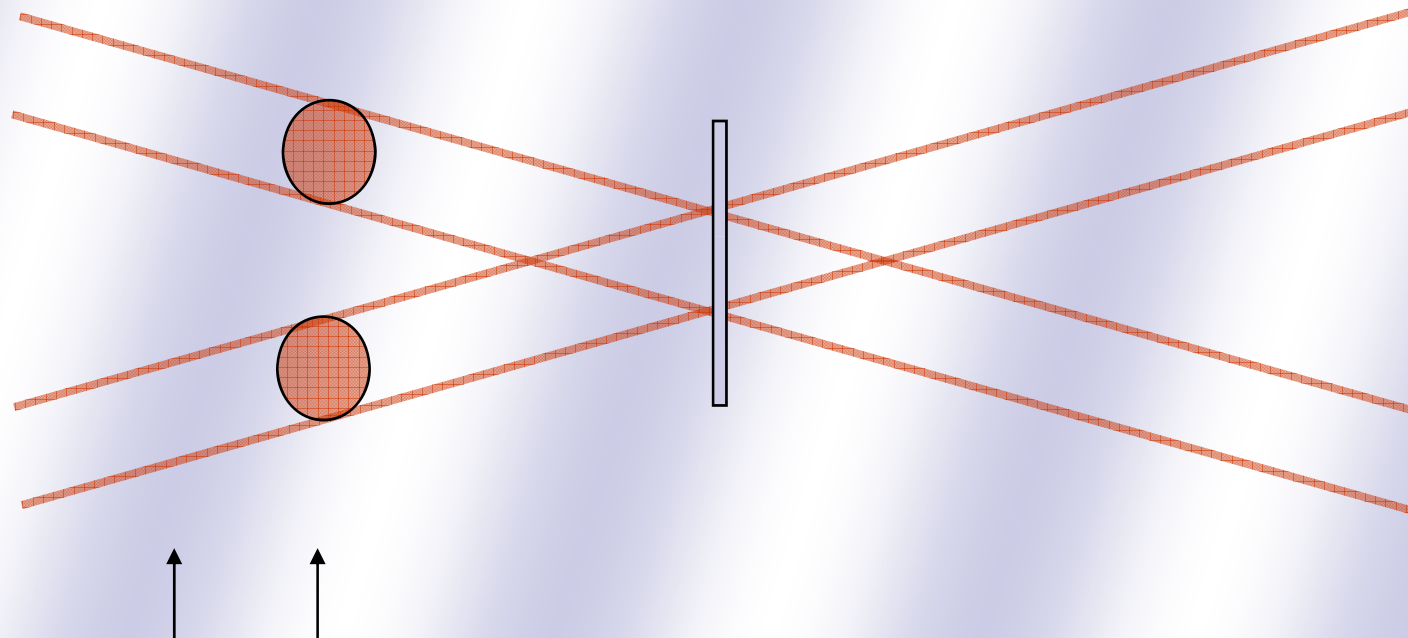




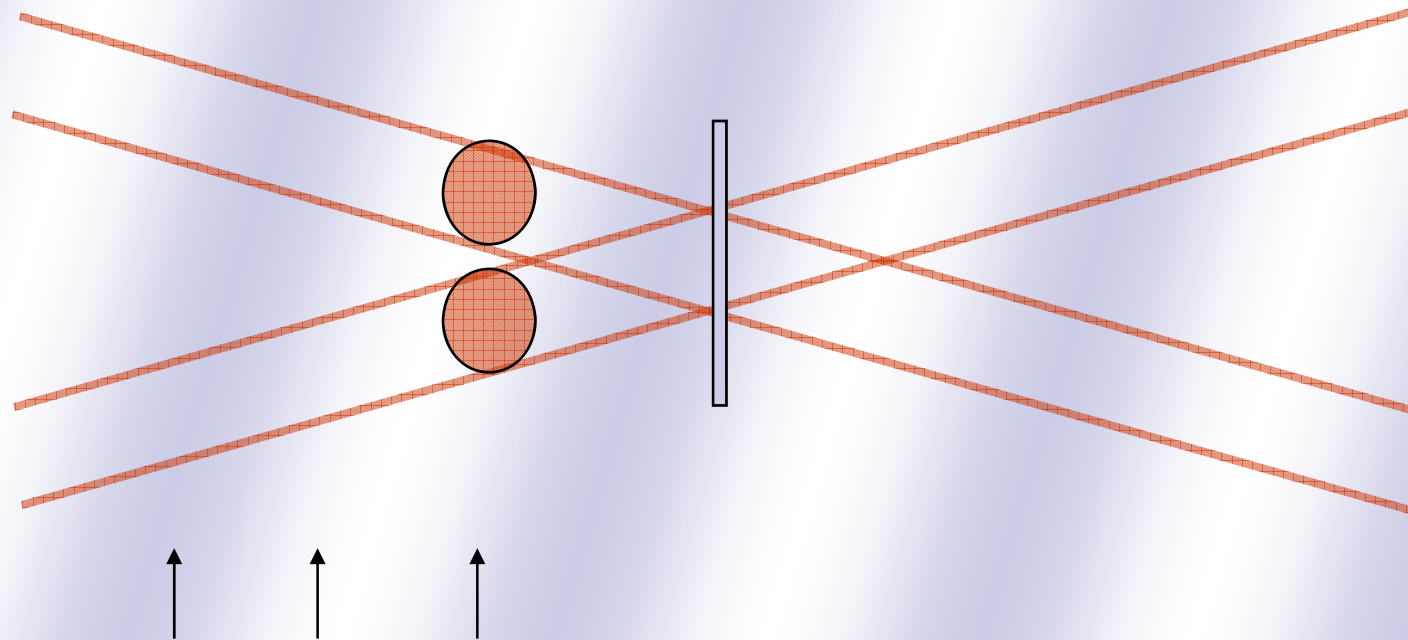
# Interference



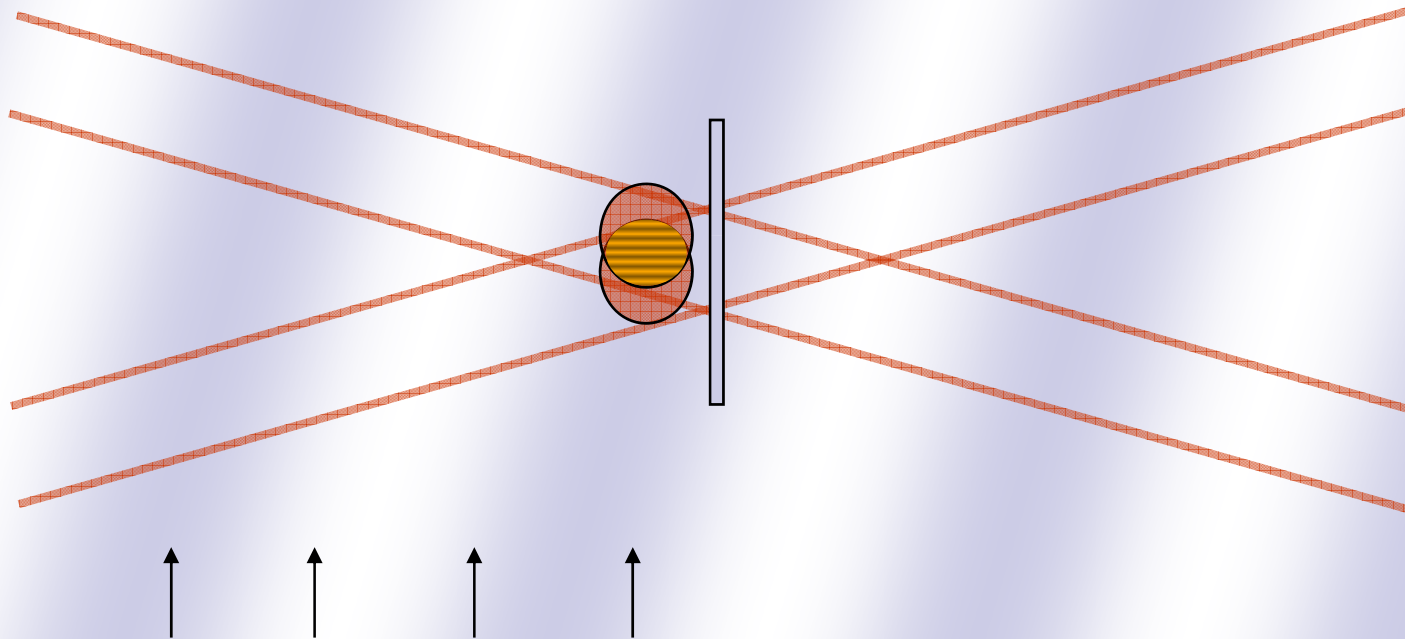
# Interference



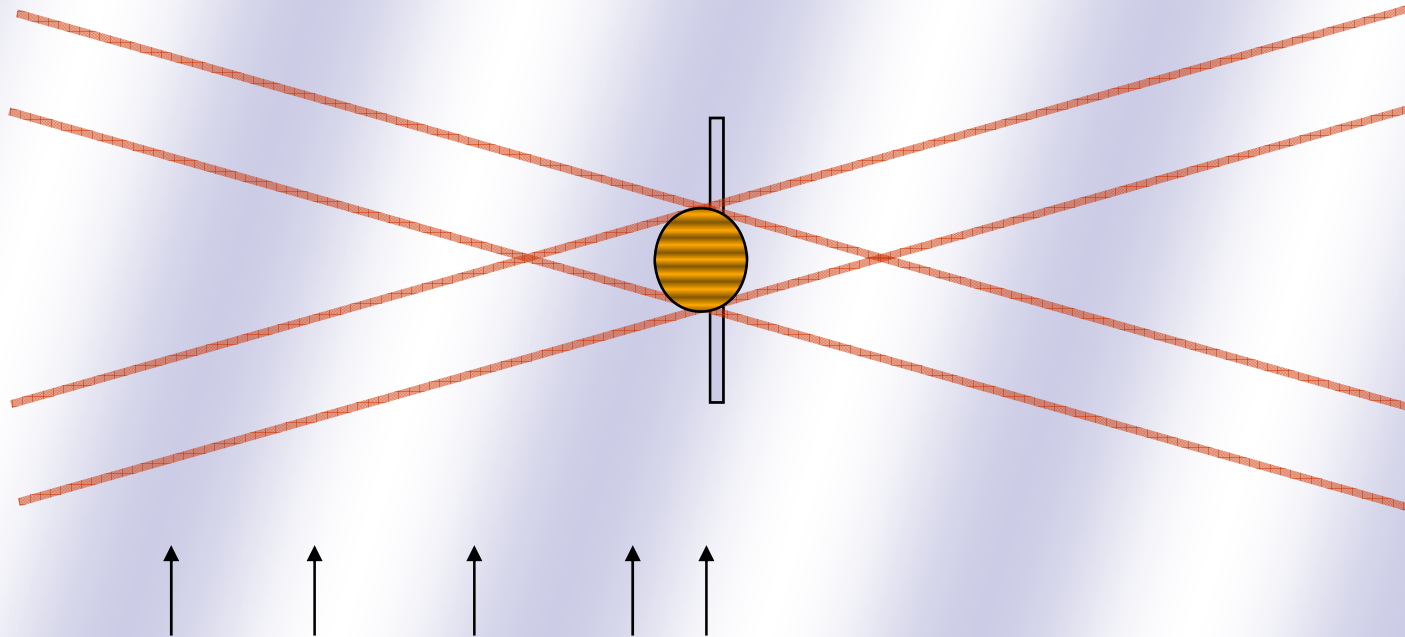
# Interference



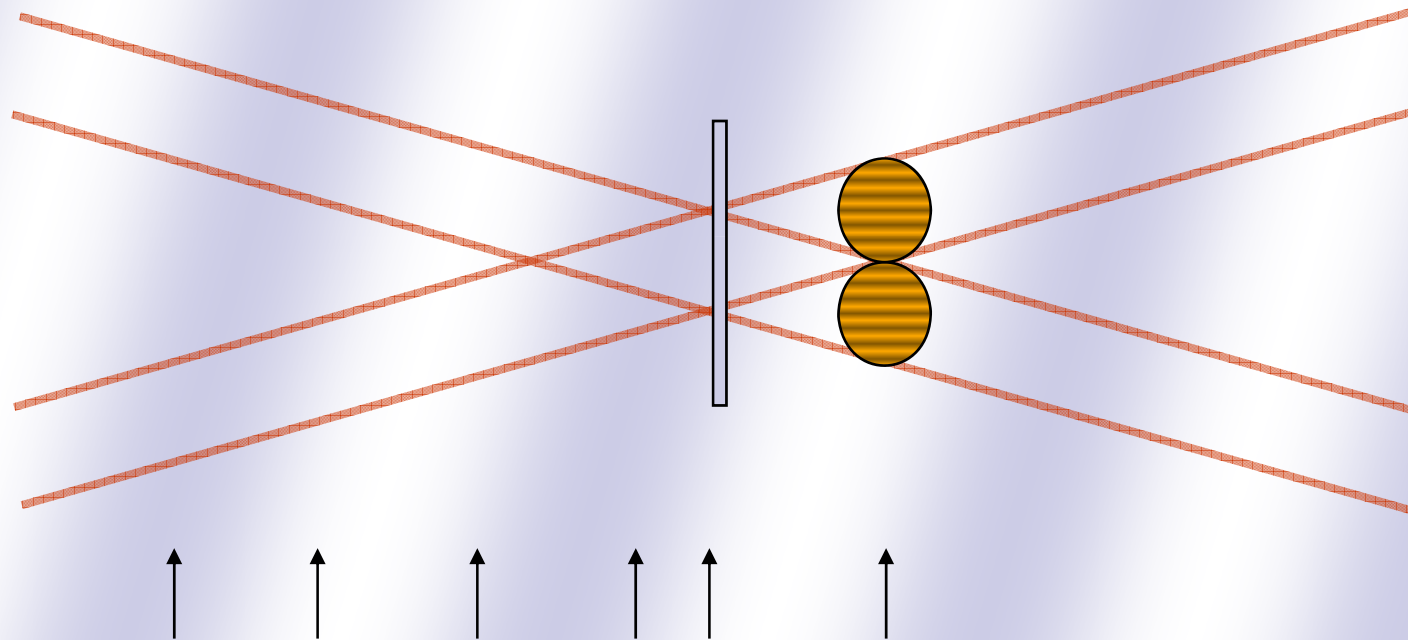
# Interference



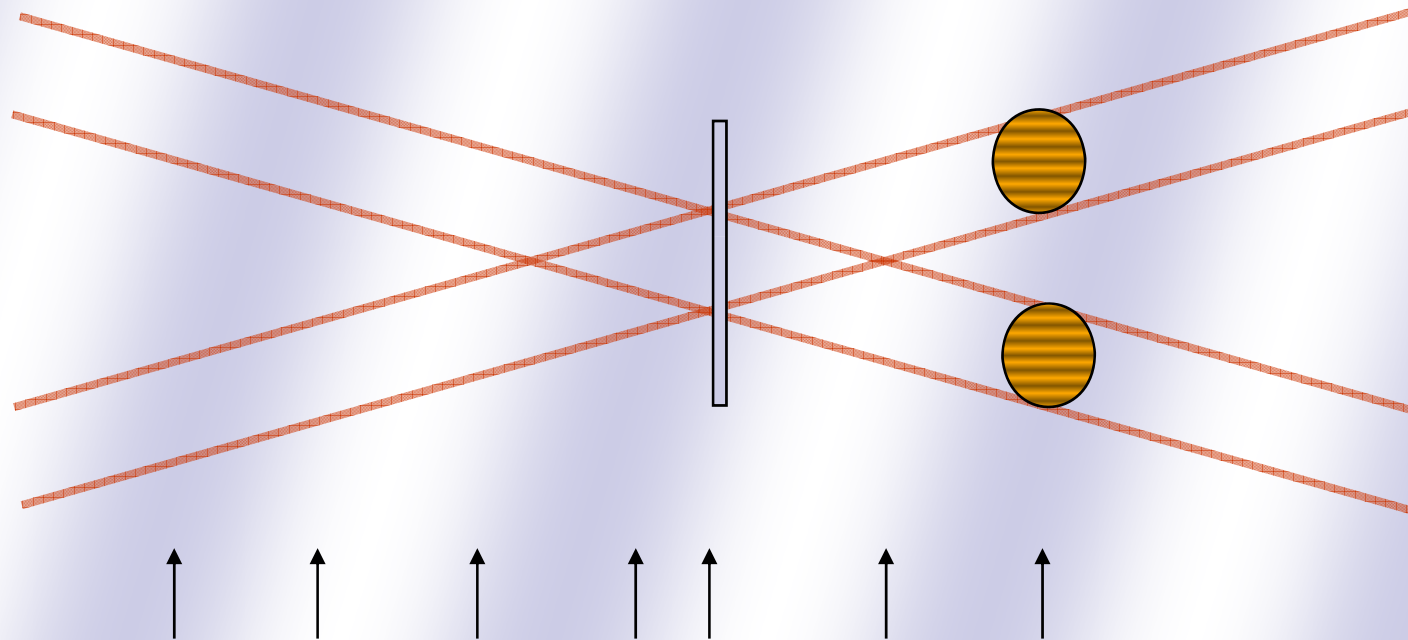
# Interference



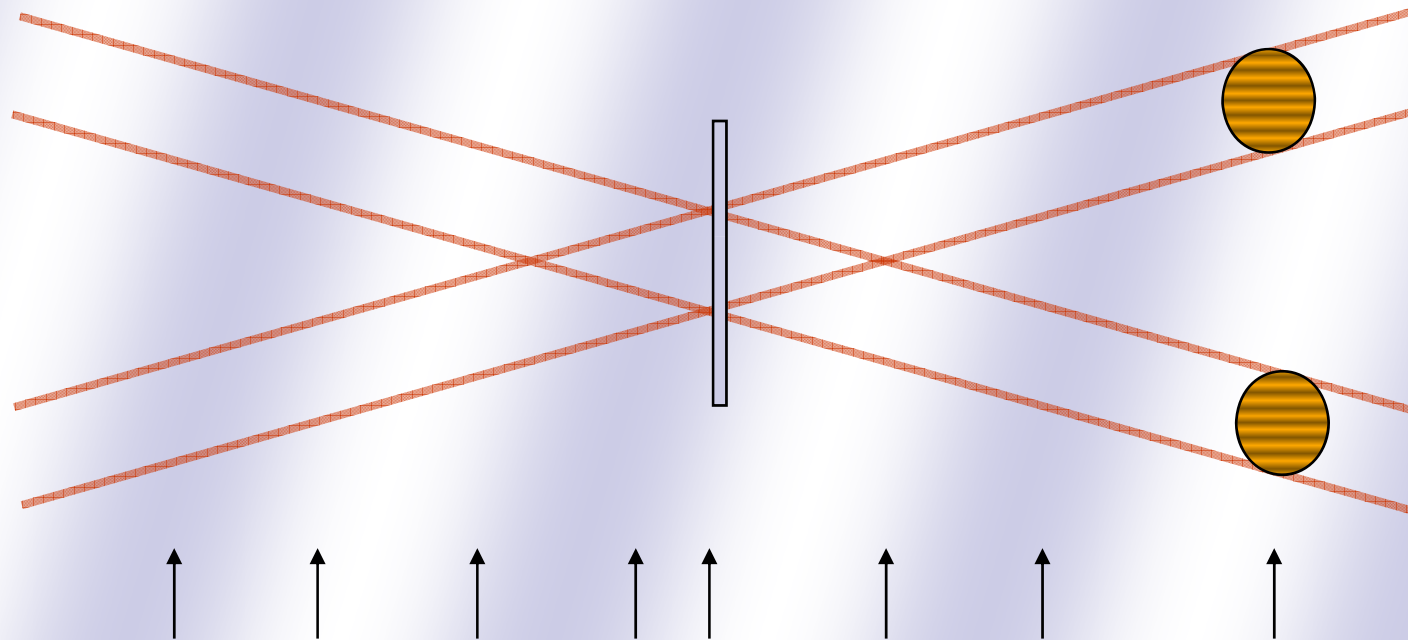
# Interference



# Interference



# Interference





# **Interference**

**There is an effective redistribution of energy due to the matter inserted in the optical path**

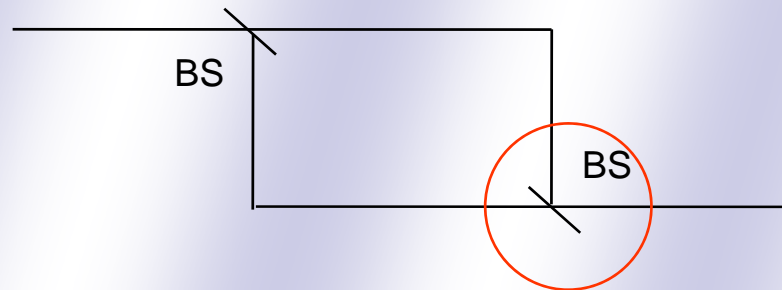
**In this case the redistribution is nonlocal**

# Interference

The fundamental step is the observation of the interference pattern:

$$I = \frac{1}{2} c \varepsilon_0 |E_1 + E_2|^2$$

Lets see it in more detail on an interferometer:

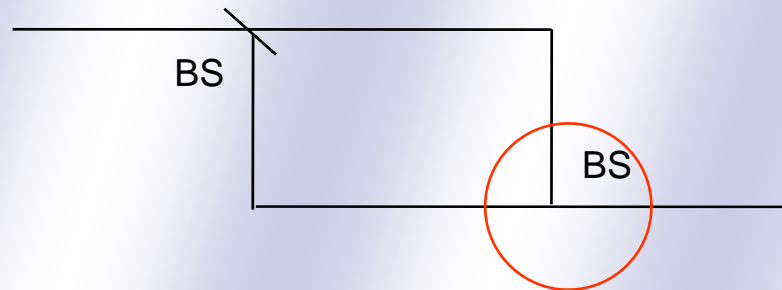


# Interference

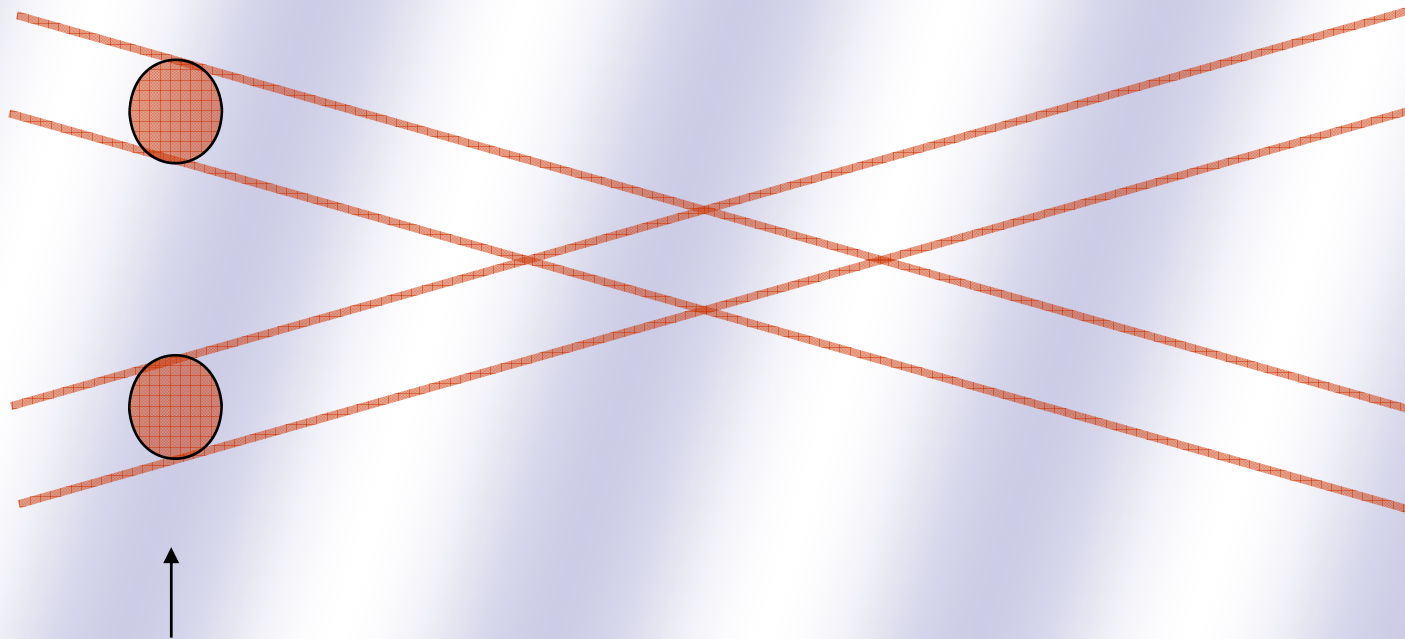
The fundamental step is the observation of the interference pattern:

$$I = \frac{1}{2} c \varepsilon_0 |E_1 + E_2|^2$$

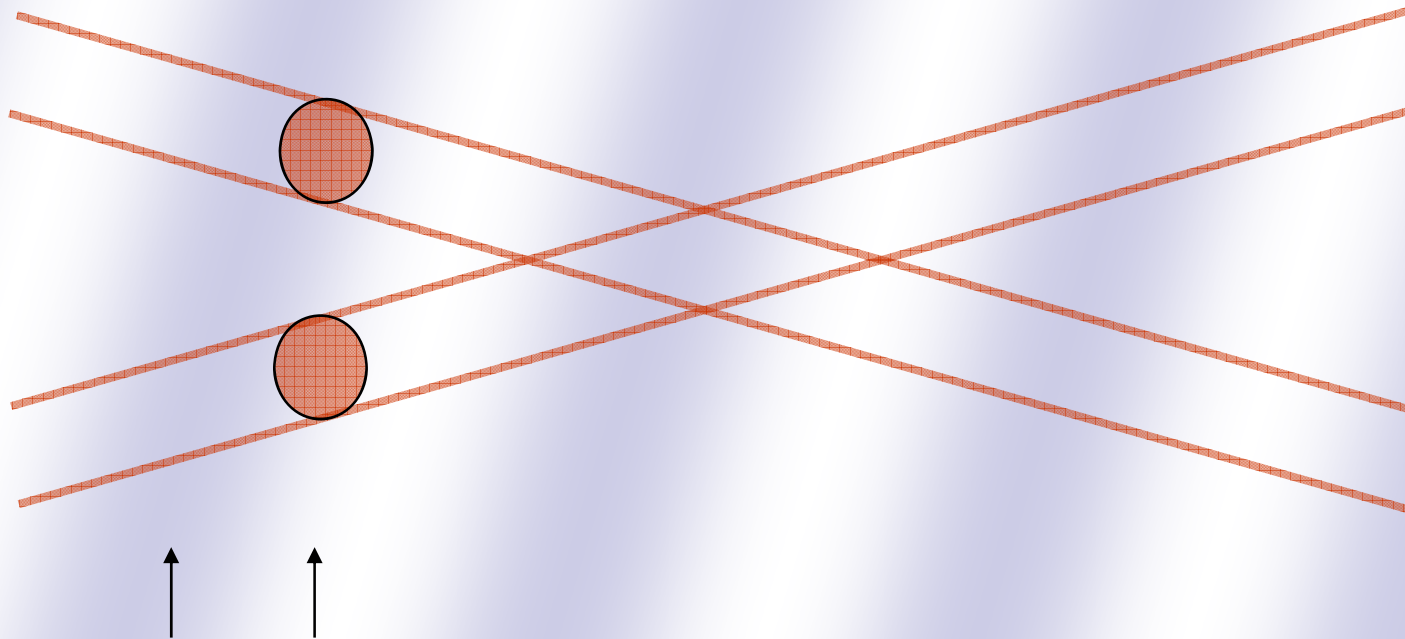
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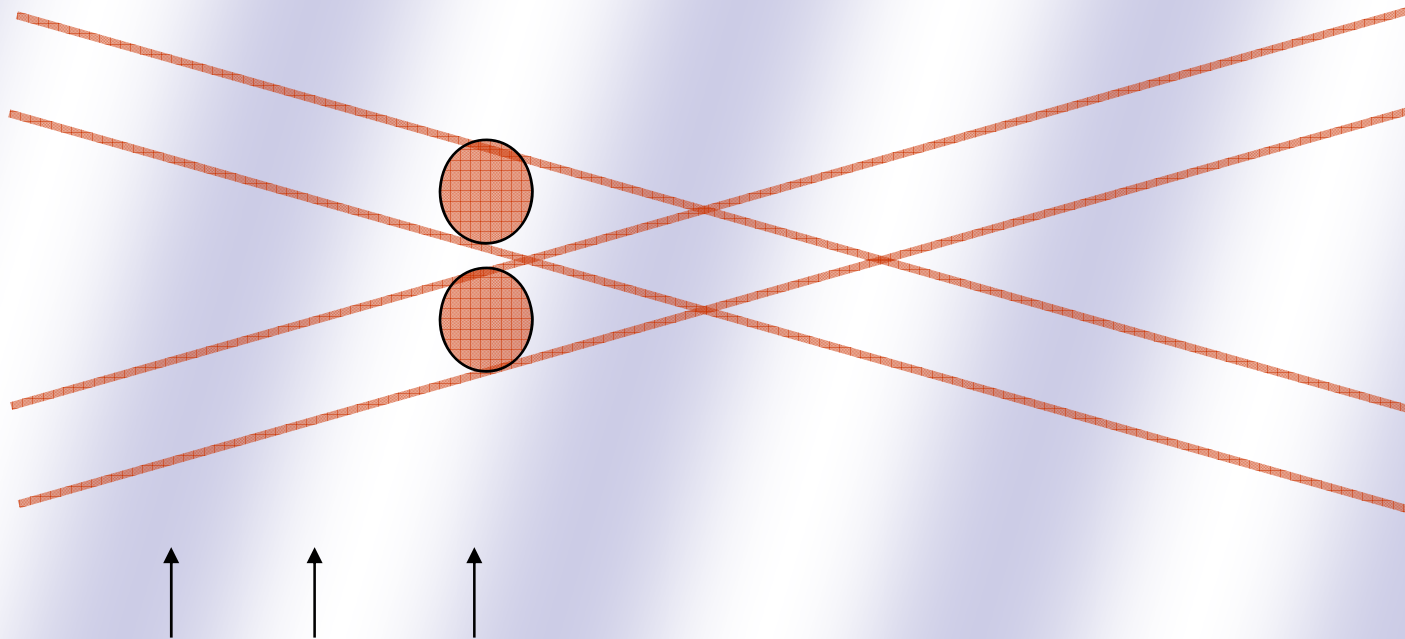
# Interference



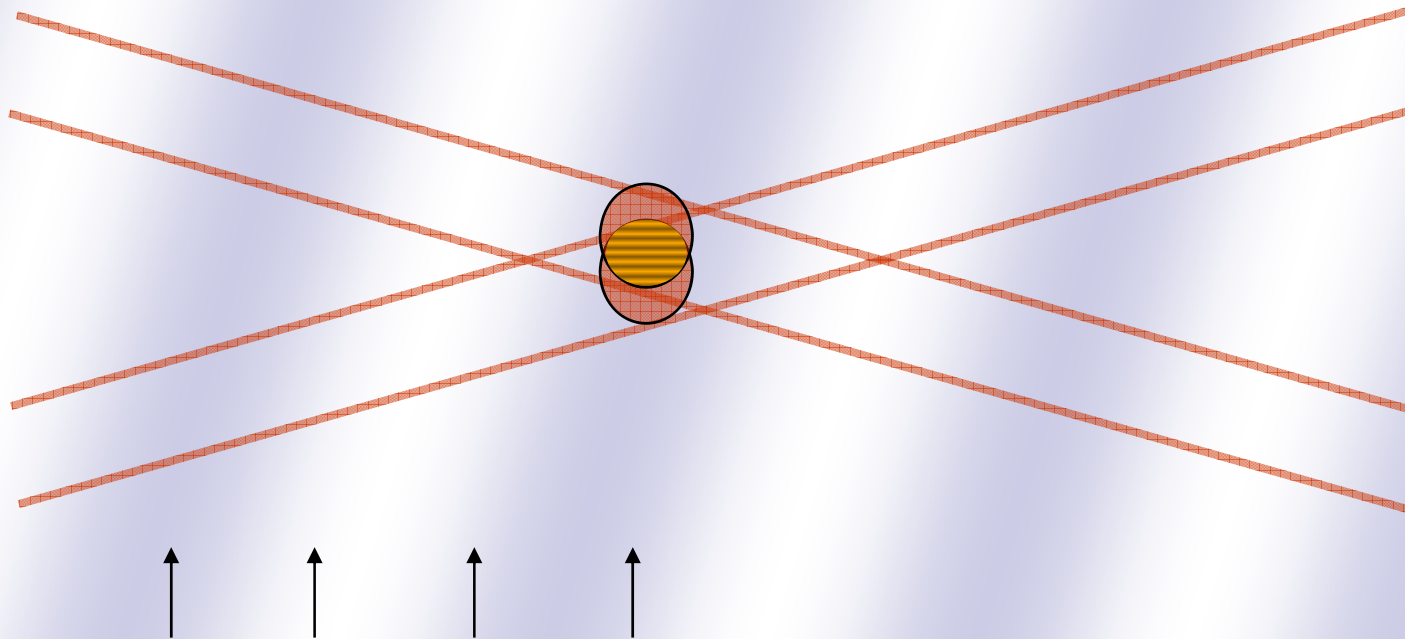
# Interference



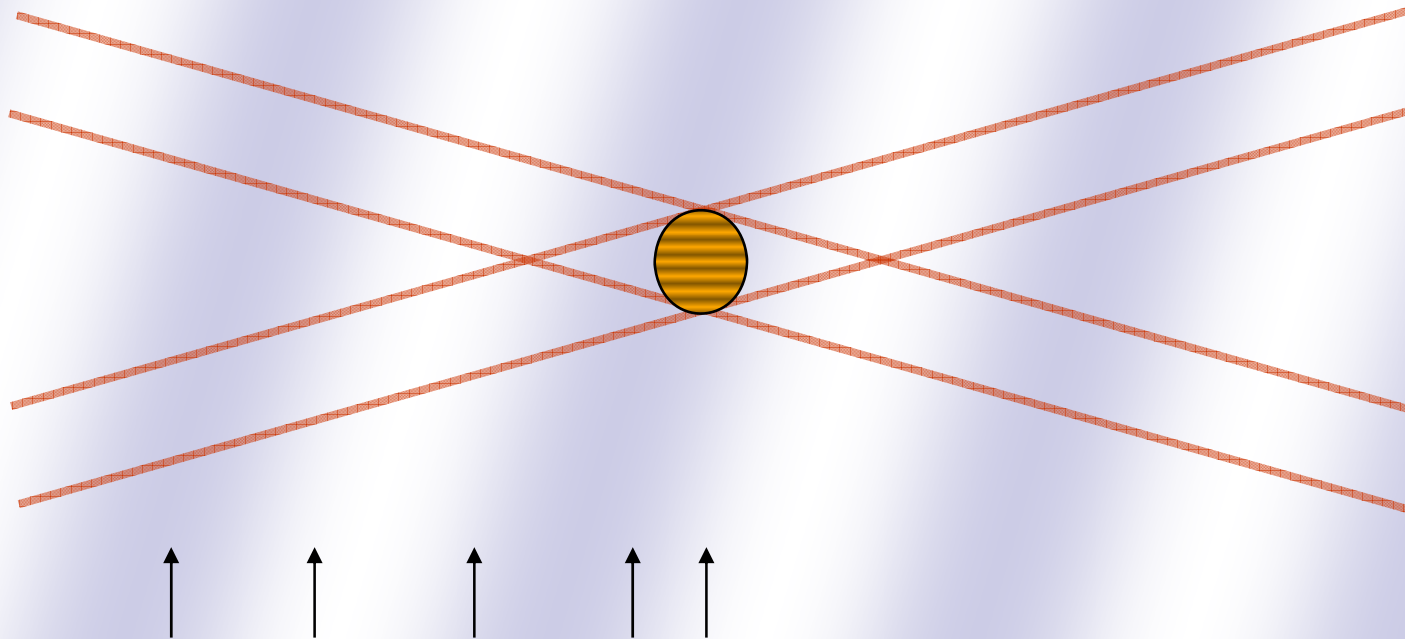
# Interference



# Interference

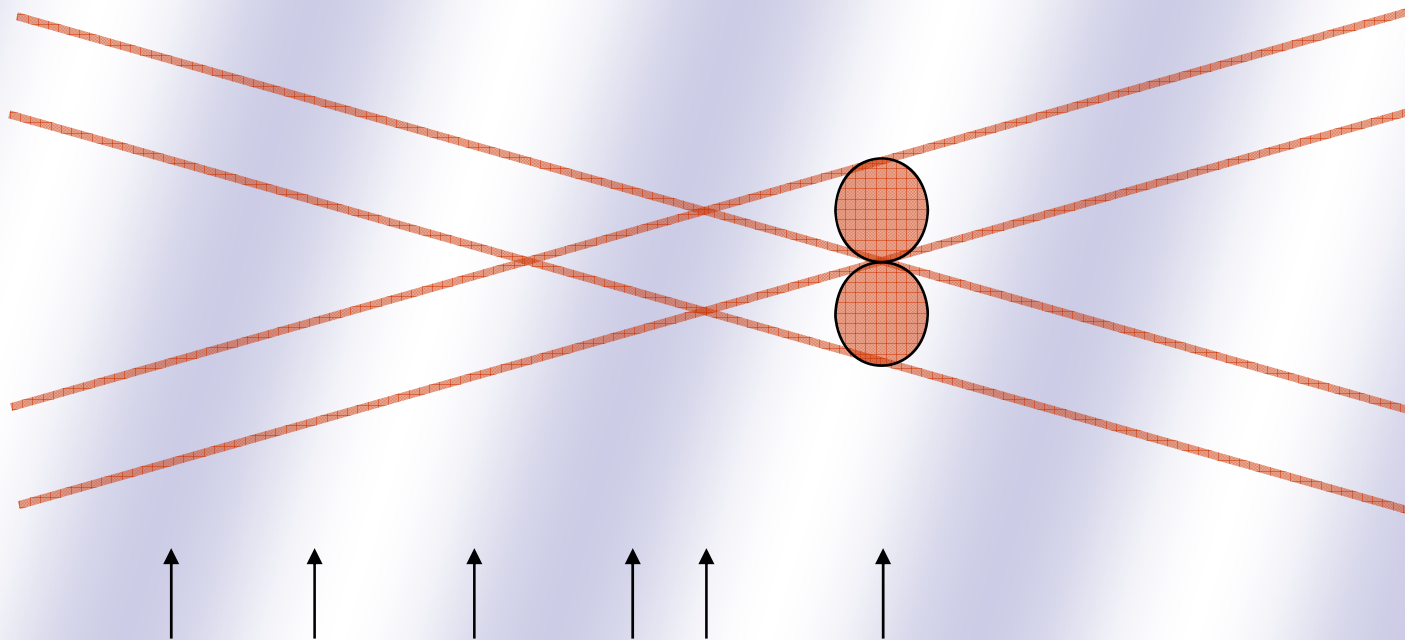


# Interference

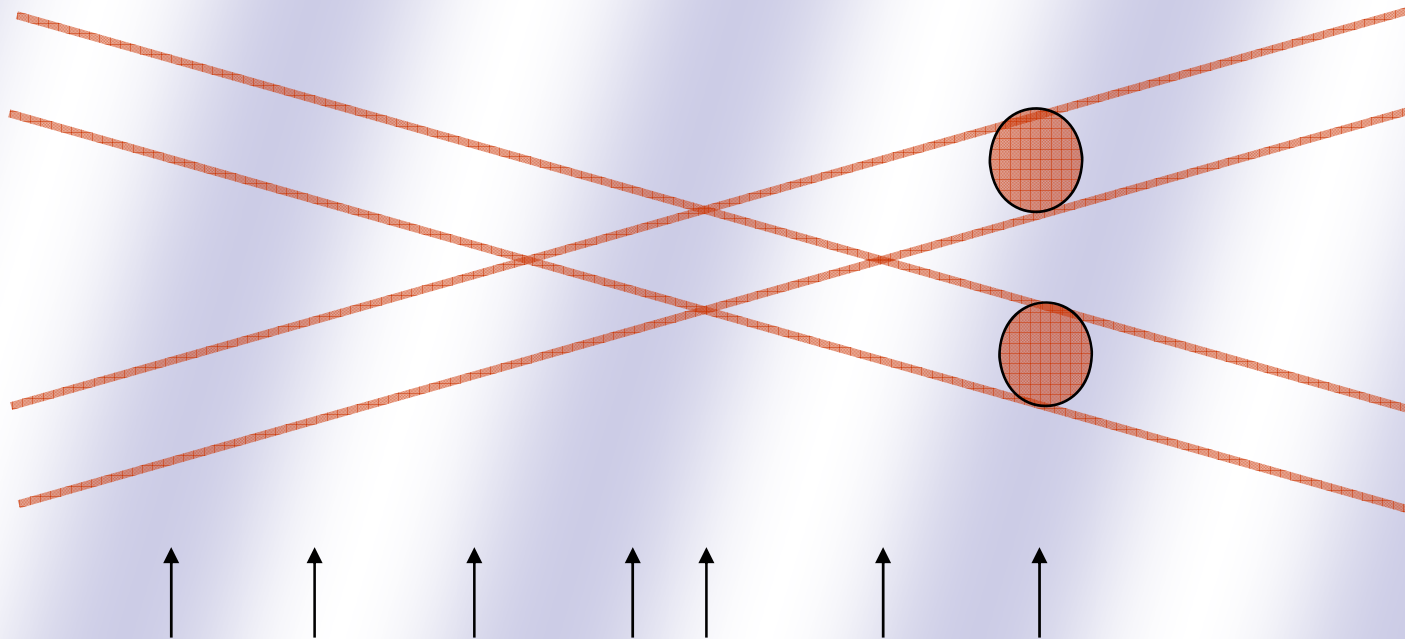




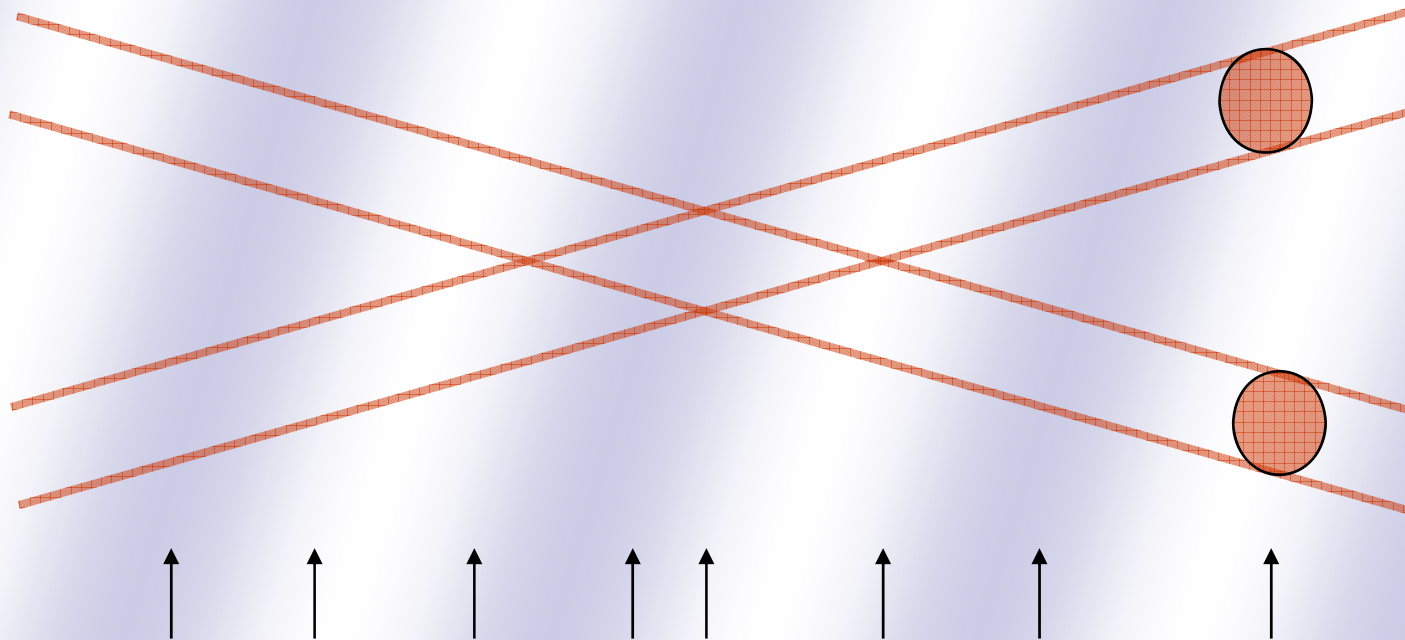
# Interference



# Interference



# Interference



# Interference

**There is an effective redistribution of energy only on the screen inserted in the optical path**

**In this case the redistribution is local**

# Interference

**Therefore there is a difference if we insert or not a combining beam-splitter**

- **Electric field superposition is common to both scenarios**
- **Interference is not**

**Before going any further lets discuss classical NLO**

# Nonlinear Optics

$$P = \varepsilon_0 \tilde{\chi} \cdot \vec{E}$$

$$\begin{aligned} \vec{D} &= \varepsilon_0 \vec{E} + \vec{P} \\ &= \varepsilon_0 (1 + \tilde{\chi}) \vec{E} \end{aligned}$$

**Electronic nonlinearities:**

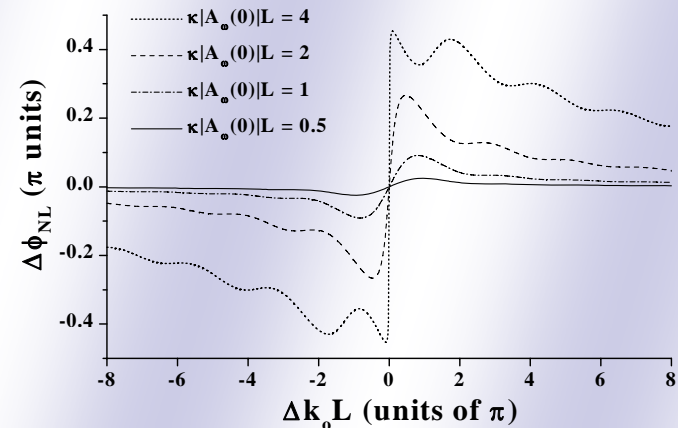
$$P = \varepsilon_0 \tilde{\chi}^{(1)} \cdot \vec{E} + \varepsilon_0 \tilde{\chi}^{(2)} : \vec{E} + \varepsilon_0 \tilde{\chi}^{(3)} \vdots \vec{E} + \dots$$

**Second order NLO  $\chi^{(2)}$ : SHG, DFG, SFG, OR, PG**

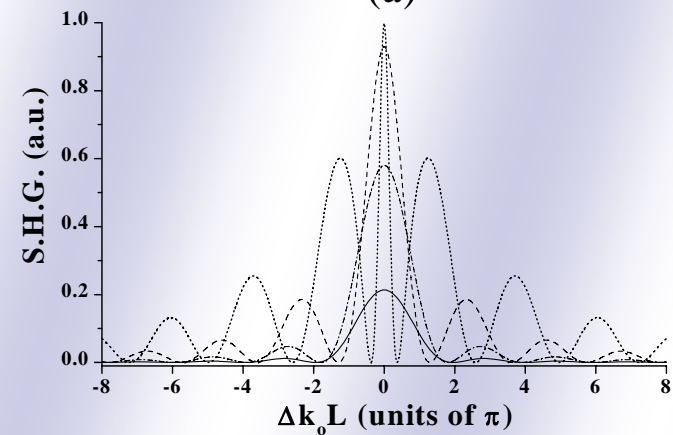
**Third order NLO  $\chi^{(3)}$ : Kerr, Raman, SPM, XPM, ....**

# Second Harmonic Generation

Second Harmonic Generation describes the interaction between an intense beam and a noncentrosymmetric material which produces a beam with higher energy (double) than the incident



(a)



(b)

# Interference Wave-Matter

- For spatially varying distribution of parameters (type I SHG)

$$\frac{dA_{\omega}(z)}{dz} + i[k_{0,\omega} + \delta k_{\omega}(z)] A_{\omega}(z) = -i \kappa(z) A_{2\omega}(z) A_{\omega}^*(z)$$

$$\frac{dA_{2\omega}(z)}{dz} + i[k_{0,2\omega} + \delta k_{2\omega}(z)] A_{2\omega}(z) = -i \kappa(z) A_{\omega}(z) A_{\omega}(z)$$

→ define: -  $\Delta k_0 = 2k_{0,\omega} - k_{0,2\omega}$  ; -  $\Delta K(z) = 2\delta k_{\omega}(z) - \delta k_{2\omega}(z)$

$$-G(z) = L \cdot U(z/L) \kappa(z) e^{i \int_0^L \Delta K(z') dz'}$$

(U(z): *Rectangle function*)

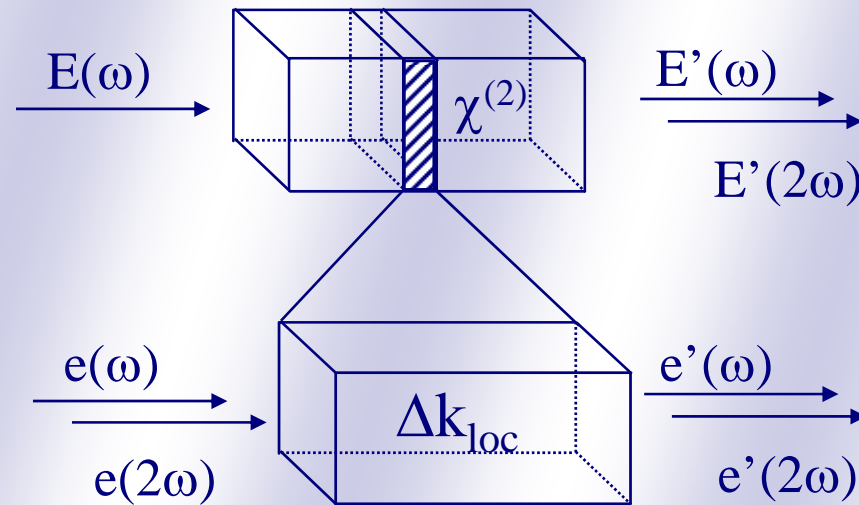
⇒

$$\frac{d\mathcal{A}_{\omega}(z)}{dz} = -i G(z) \mathcal{A}_{2\omega}(z) \mathcal{A}_{\omega}^*(z) e^{i \Delta k_0 z}$$

$$\frac{d\mathcal{A}_{2\omega}(z)}{dz} = -i G(z) \mathcal{A}_{\omega}(z) \mathcal{A}_{\omega}(z) e^{-i \Delta k_0 z}$$



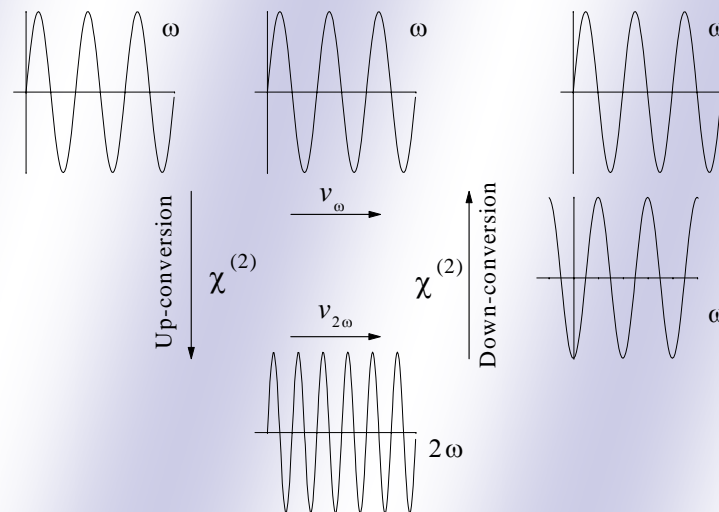
# Second Harmonic Generation



- SHG is a succession of local seeding processes (a non-local process).
- Local SHG evolution along propagation depends on:
  - Amplitude and Phase of Transmitted Fundamental ( $\omega$ ).
  - Amplitude and Phase of Generated SH ( $2\omega$ ).
  - Local wave-vector mismatch ( $\Delta k_{loc}$ ) and nonlinear coupling coefficient.

# Second Harmonic Generation

The SHG process is governed by phase-matching conditions to achieve a larger or smaller conversion efficiencies

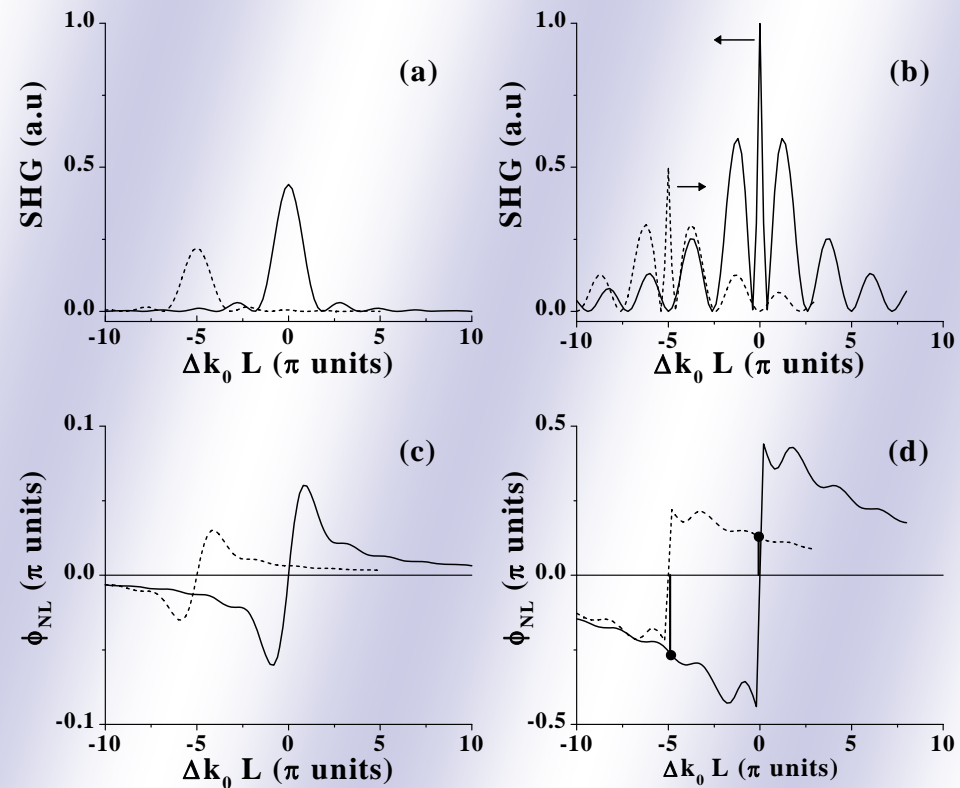


**Cascading - interference**

# Second Harmonic Generation

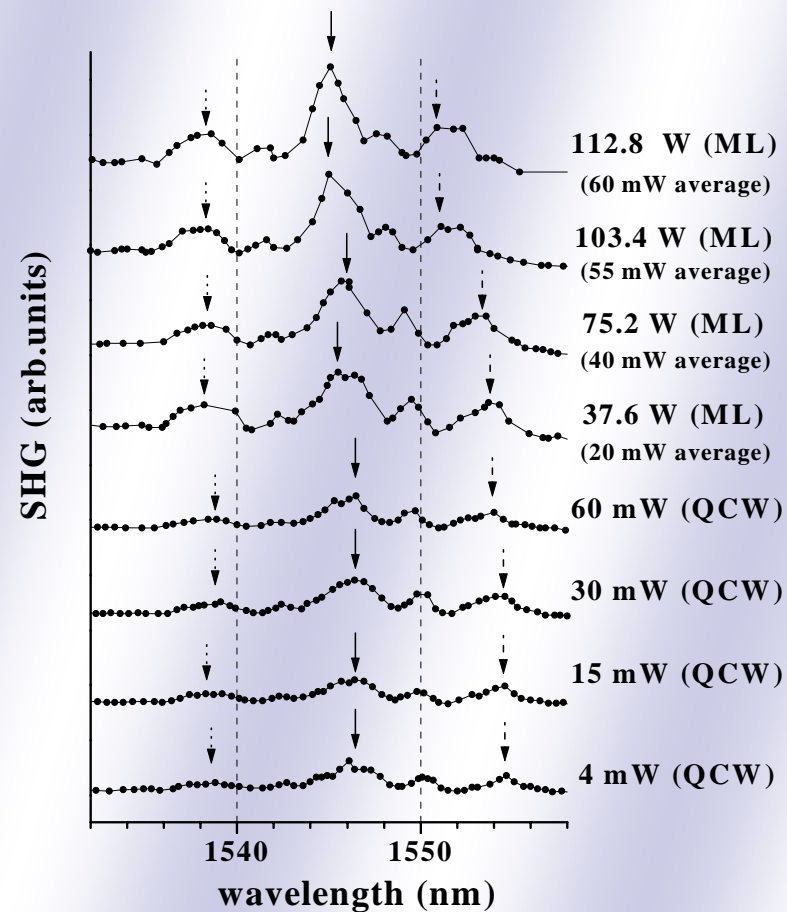
**This is more evident if there are two or more simultaneous “resonances” in which SHG can be achieved**

Cascading - multimode



# Second Harmonic Generation

**This is more evident if there are two or more simultaneous “resonances” in which SHG can be achieved**



# Interference Wave-Matter

- For spatially varying distribution of parameters (type I SHG)

$$\frac{dA_{\omega}(z)}{dz} + i[k_{0,\omega} + \delta k_{\omega}(z)] A_{\omega}(z) = -i \kappa(z) A_{2\omega}(z) A_{\omega}^*(z)$$

$$\frac{dA_{2\omega}(z)}{dz} + i[k_{0,2\omega} + \delta k_{2\omega}(z)] A_{2\omega}(z) = -i \kappa(z) A_{\omega}(z) A_{\omega}(z)$$

→ define: -  $\Delta k_0 = 2k_{0,\omega} - k_{0,2\omega}$  ; -  $\Delta K(z) = 2\delta k_{\omega}(z) - \delta k_{2\omega}(z)$

$$-G(z) = L \cdot U(z/L) \kappa(z) e^{i \int_0^L \Delta K(z') dz'}$$

(U(z): *Rectangle function*)

⇒

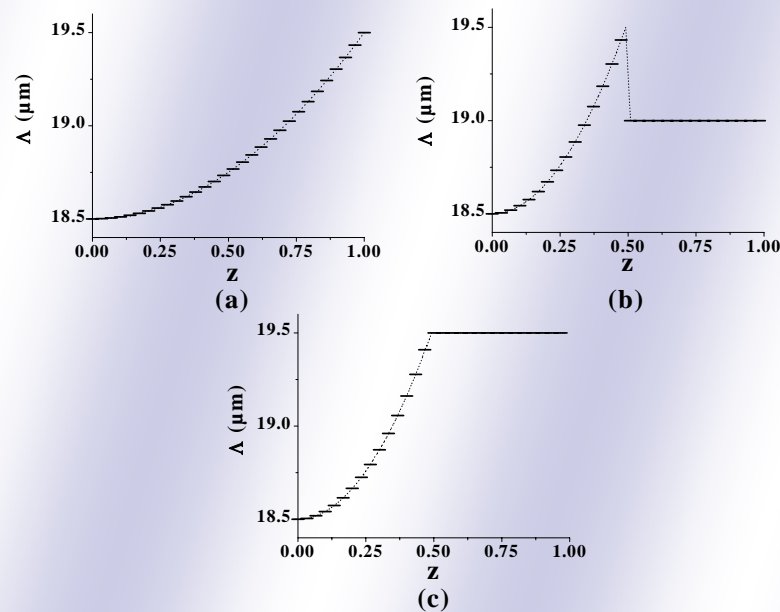
$$\frac{d\mathcal{A}_{\omega}(z)}{dz} = -i G(z) \mathcal{A}_{2\omega}(z) \mathcal{A}_{\omega}^*(z) e^{i \Delta k_0 z}$$

$$\frac{d\mathcal{A}_{2\omega}(z)}{dz} = -i G(z) \mathcal{A}_{\omega}(z) \mathcal{A}_{\omega}(z) e^{-i \Delta k_0 z}$$

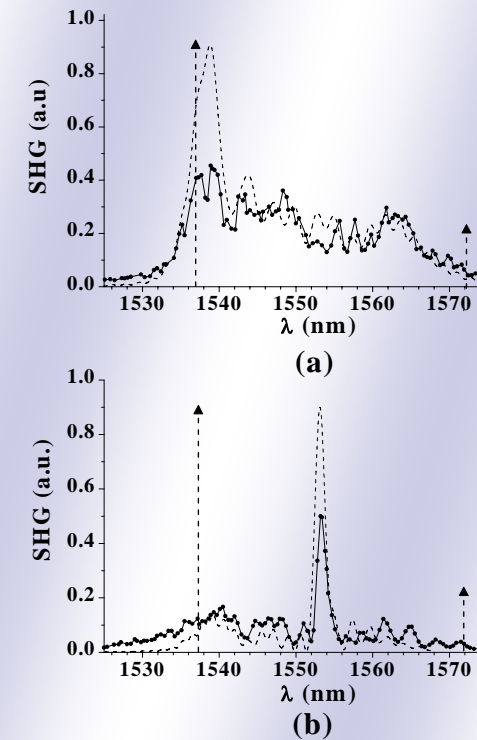
# Interference Wave-Matter

The response can be rather complex, but predictable

## Nonuniform profiles

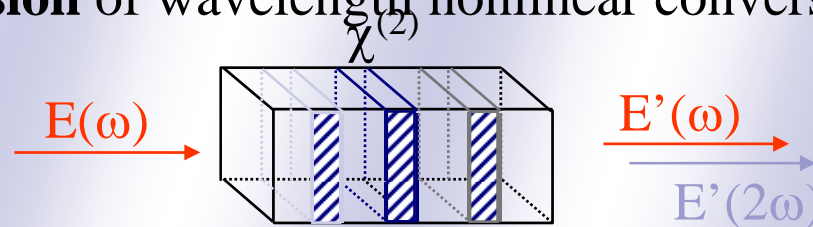


## SHG detuning curves

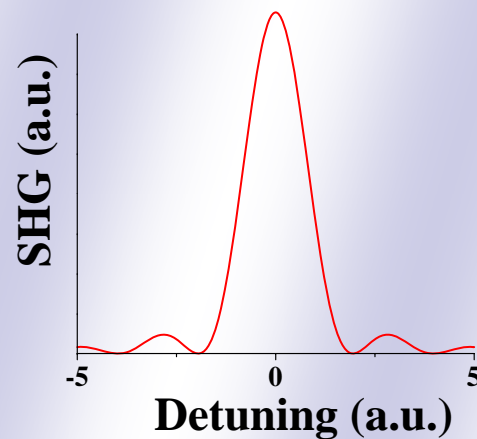


# Interference Wave-Matter

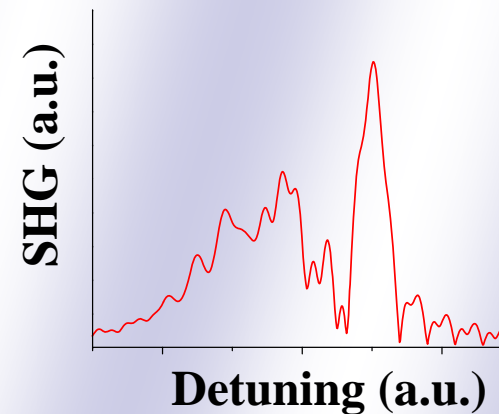
- SHG is a **succession** of wavelength nonlinear conversion process (*nonlocal*)



- The collective behaviour is observed in the SHG detuning curve



Uniform  $\rightarrow$  **sinc-like**



Nonuniform  $\rightarrow$  **complex**

**This is nothing more than standard interferometry in which the medium modulates the fringe “visibility”**



# Comparing Interferometry and NLO

**Interferometry:** the light redistribution is used to describe what happened to the light as it is modified by a dense medium, after it has been observed of course – in coordinate space

**NLO:** the light redistribution is used to describe what happened to the dense medium as it is modified by light, after it has been observed of course – in wavevector space



# Comparing Interferometry and NLO

**This effect is not particular to interferometry-  
second order NLO**

**Both phenomena discussed are related by  
swapping coordinate and wave-vectors.**

**Similar equivalences can be found between  
frequency and time (say OCT and local NLO)**

**This is to be expected because the way these  
spaces are included in the EM phenomenology**

$$\vec{E} = \vec{E}_0 e^{i(\vec{k} \cdot \vec{r} + \omega t + \phi)}$$



# Conclusions

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# Question

**1. Is EM interference a basic phenomena or does it require a dense medium to exist?**

*A: According to me, interference is NOT a basic phenomena*

*EM field superposition sampled by radiation-matter interaction IS.*