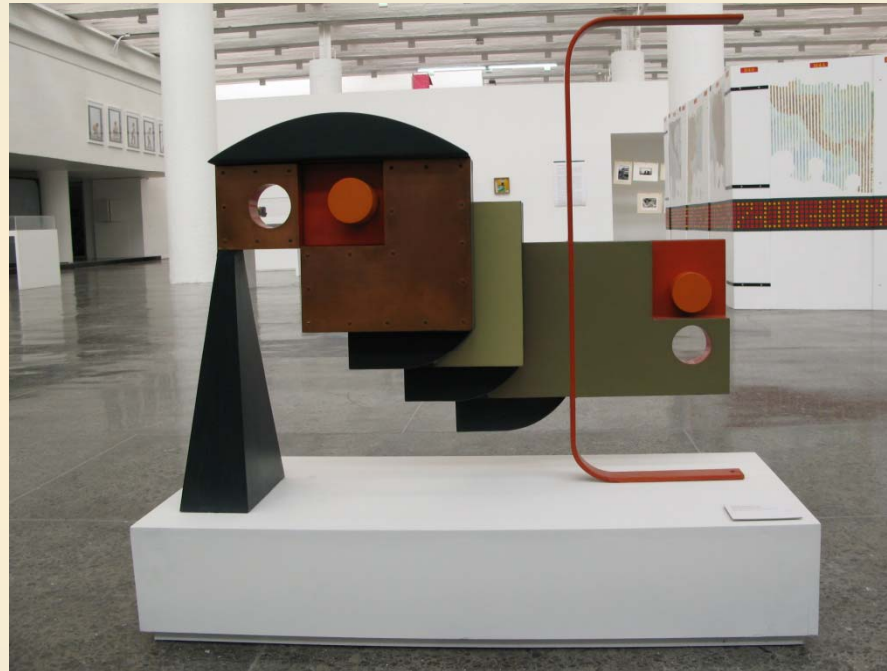




Zero-point energy: the idea underlying the quantum



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● Introduction

T. H. Boyer, Am. J. Phys. **71**, 866 (2003):

By a thermodynamic analysis of the harmonic oscillator Boyer finds:

$$\begin{array}{ccc}
 U \approx k_B T & \xrightarrow{T=0} & U \approx \frac{1}{2} \hbar \omega + k_B T \\
 \text{Wien's law} & & \text{Zero-point energy} \\
 & & \text{Breakdown of equipartition!!} \\
 & & U \approx k_B T
 \end{array}$$

The existence of a zero-point energy emerges **naturally** from the thermodynamic relations and breaks down with classical physics...

How far can we get by introducing this new (non-classical) physical element into the statistical description of a system of harmonic oscillators?

The results obtained coincide with those that led Planck and Einstein to establish the quantum hypothesis of the radiation field.





OUTLINE

- **Thermodynamic and statistical description of the harmonic oscillator**

Introduction the zero-point energy into the statistical description

- **Establishing Planck's law**

Derivation of Planck's law from the statistical information

- **Discussion: the quantum hypothesis of Planck and Einstein**

Historical overview and comparison of Einstein's analysis with the present approach

- **Continuous vs discrete**

Disclosing the origin of discreteness

- **Non-thermal fluctuations**

Introduction of the zero-point energy fluctuations into the description

- **Concluding remarks**



● Thermodynamic and statistical description of the harmonic oscillator

- Statistical information

The statistical information of the system is contained in a distribution function W from which the mean values of any function $f(E)$ can be calculated:

$$\overline{f(E)} = \int f(E) W(E) dE \quad \overline{E} = \int U W(E) dE$$

Moreover, W is such that it maximizes the entropy defined as

$$S = -k \int W \ln W dE$$

Such distribution can be expressed as (Montroll and Shlesinger, Jour. Stat. Phys. **32**, 209 (1983))

$$W_g(E) = \frac{1}{Z_g} g(E) e^{-\beta E}$$

with

$$Z_g = \int g(E) e^{-\beta E} dE \quad \beta = \frac{1}{kT}$$



• Thermodynamic and statistical description of the harmonic oscillator

The choice $g(E)=1$ corresponds to $W_g = \frac{1}{Z_1} e^{-\beta E}$, and leads to

$$\bar{E} = U = kT \quad \text{Equipartition and inconsistency with some of Boyer's relations}$$

Thus in order to account for a zpe a more general function $g(E)$ is needed.

• Thermal fluctuations

For the general distribution W_g the following recurrence relation holds (r a positive integer)

$$\overline{E^{r+1}} = U \overline{E^r} + \frac{d\overline{E^r}}{d\beta}$$

In particular, for the dispersion of the energy,

$$\sigma_E^2 = \overline{E^2} - U^2$$

we obtain

$$\sigma_E^2 = \frac{dU}{d\beta} = kT^2 C_{\eta} \xrightarrow{T=0} \sigma_E^2 = 0$$

fluctuations are of thermal nature



- Establishing Planck's law

Knowing ρ_{ω}^2 as a function of U determines $U(\omega)$

$$d\rho_{\omega}^2 = \frac{dU}{\rho_{\omega}^2}, \quad \rho_{\omega}^2 \rho_{\omega_0}^2 = 0$$

- Determination of the dispersion as a function of U

$$\rho_{\omega}^2 * E^2 = \rho_{\omega_0}^2$$

We expand ρ_{ω}^2 as a power series of U and note that for high temperatures the result $\rho_{\omega}^2 \propto U^2$ must be recovered. Thus the dispersion reduces to a polynomial of second order in U :

$$\rho_{\omega}^2 \rho_{\omega_0}^2 = a_0 + a_1 U + U^2$$

Since ρ_{ω}^2 is invariant under the transformation $E \rightarrow \lambda E$ it follows that $\rho_{\omega}^2 \rho_{\omega_0}^2 = \rho_{\omega}^2 \rho_{\omega_0}^2$, thus

$$\rho_{\omega}^2 \rho_{\omega_0}^2 = U^2$$

And since $\rho_{\omega}^2 \rho_{\omega_0}^2 = 0$ it follows that

$$\rho_{\omega}^2 \rho_{\omega_0}^2 = U^2 \neq E_0^2$$



Establishing Planck's law

The equation

$$\frac{dU}{U^2 - E_0^2} = \frac{dE_0}{E_0^2}$$

has the solution

$$\left\{ \begin{array}{ll} \frac{1}{U} & \text{for } E_0 = 0 \rightarrow U = \frac{1}{E_0} \quad \text{No zpe means equipartition} \\ \frac{1}{E_0} \coth^{-1} \frac{U}{E_0} & \text{for } E_0 \neq 0 \rightarrow U = E_0 \coth E_0 \quad \text{Zpe means Planck's law} \end{array} \right.$$

The zero-point energy is a **necessary and sufficient** condition to arrive at Planck's law...

(no quantum hypothesis is required)

$$U = U_T - E_0 \rightarrow U_T = E_0 \coth E_0 \quad \text{Original result of Planck}$$



● Discussion: the quantum hypothesis of Planck and Einstein

$$\frac{\frac{2S}{U_T^2}}{\frac{2S}{U_T^2}} = \frac{1}{T}$$

● Planck's analysis: the birth of the quantum

$$\left. \begin{array}{l} T \text{ (clock icon)} \\ \frac{2S}{U_T^2} \text{ (hand icon)} \end{array} \right\} \frac{k}{U_T^2}$$

$$\left. \begin{array}{l} T \text{ (clock icon)} \\ U_T \text{ (hand icon)} \\ 2E_0 e^{2E_0/T} \text{ (hand icon)} \end{array} \right\} \frac{k}{2E_0 U_T}$$

$$\frac{2S}{U_T^2} \text{ (hand icon)} \left\{ \frac{k}{U_T^2 \cdot 2E_0 U_T} \right.$$

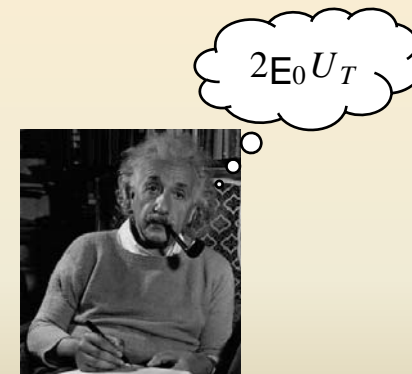


● Einstein's analysis: the birth of the quantized radiation field

$$\left. \begin{array}{l} \frac{2S}{U_T^2} \text{ (hand icon)} \\ \frac{2S}{U_T^2} \text{ (hand icon)} \\ kT^2 C \nu_0 \text{ (hand icon)} \end{array} \right\} \frac{k}{U_T^2 \cdot 2E_0 U_T}$$

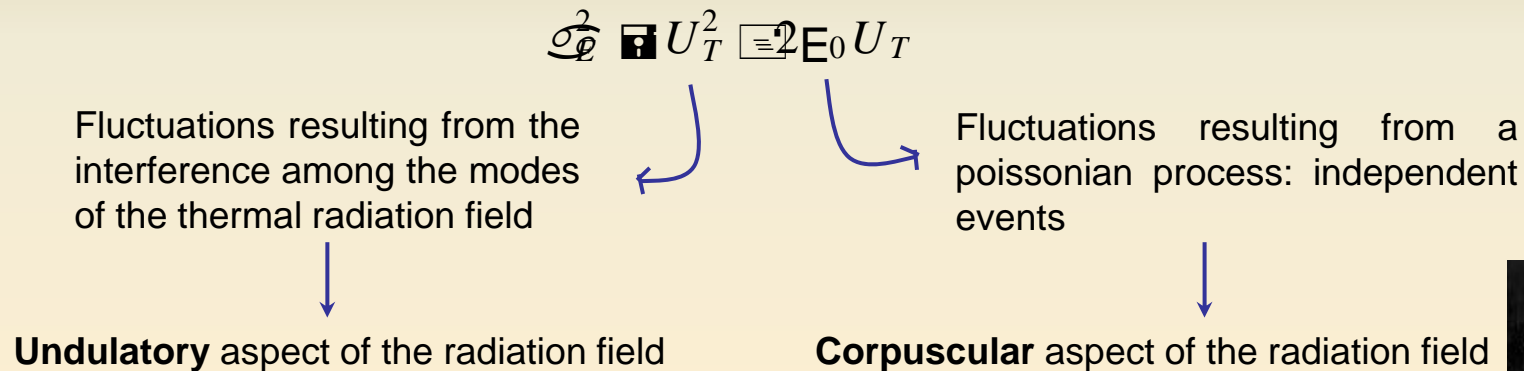
$$\frac{2S}{U_T^2} \text{ (hand icon)} \left\{ \frac{1}{U_T} \text{ (hand icon)} \left\{ \frac{1}{T^2 C \nu_0} \text{ (hand icon)} \right. \right.$$

$$kT^2 C \nu_0 \text{ (hand icon)} \left\{ \frac{U_T}{T^2 C \nu_0} \text{ (hand icon)} \left\{ \frac{2S}{U_T^2} \text{ (hand icon)} \right. \right.$$





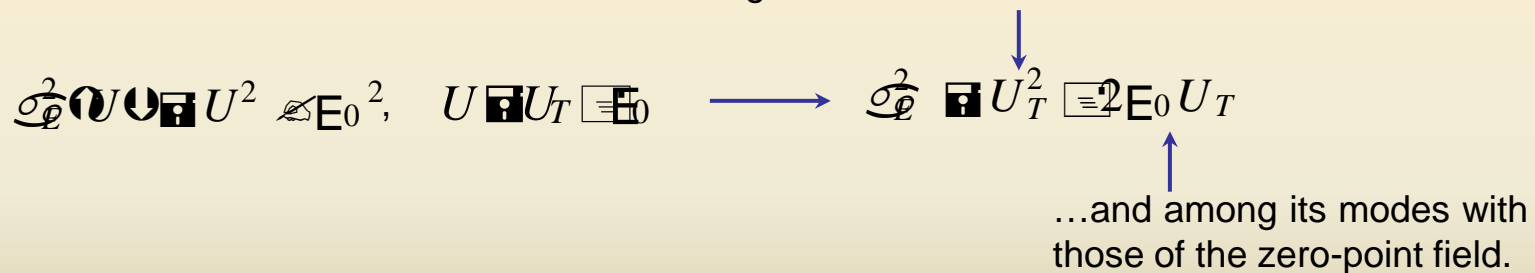
● Discussion: the quantum hypothesis of Planck and Einstein



No zpe is considered and a corpuscular property of the radiation field emerges

The present approach suggest an alternate interpretation: the fluctuations are due to the interferences of the thermal radiation field...

...among its own modes...





Continuous vs discrete

Where is the connection between both approaches?

From the thermodynamic relations it can be shown that

$$U = - \frac{d \ln Z_g}{d \beta}$$

With $U = \sum_n E_n \Omega_n$ we obtain for the partition function:

$$Z_g = \sum_n \Omega_n e^{-\beta E_n} = \sum_n \Omega_n e^{-\beta E_n} = \sum_n \Omega_n e^{-\beta E_n}$$

The partition function is related to $g(E)$ by means of the expression

$$Z_g = \int_0^\infty g(E) e^{-\beta E} dE$$

From which it follows that

$$g(E) = \sum_n \Omega_n \delta(E - E_n)$$



● Continuous vs discrete

Substituting $g(E)$ in the expression for the distribution function $W_g = \frac{1}{Z} g(E)$ leads to

$$W_g = \frac{1}{Z} \sum_n e^{-\beta E_n} \quad \text{where} \quad Z = \sum_n e^{-\beta E_n}$$

Therefore the mean value of any function $f(E)$ takes the form

$$\overline{f(E)} = \int W_g(E) f(E) dE = \frac{1}{Z} \sum_n e^{-\beta E_n} f(E_n) = \sum_n w_n f(E_n) \quad w_n = \frac{1}{Z} e^{-\beta E_n}$$

Continuous description
Averages over the **continuous variable E**

Discrete description
Averages over the **discrete states n**

Since n characterizes the energy,

$$E_n \approx \epsilon_n \epsilon_0,$$

it seems natural to assign a discrete property to the energy though it may acquire continuous values (it fluctuates!). Those mean energies corresponding to the thermal equilibrium have a discrete spectrum.



● Non-thermal fluctuations

- Including zero-point fluctuations

$$\langle \mathcal{E}^2 \rangle_{Thermal} = \langle \mathcal{E}^2 \rangle_{Thermal} + \langle \mathcal{E}_0^2 \rangle_{Thermal}, \quad \langle \mathcal{E} \rangle_{Thermal} = \langle \mathcal{E}_0 \rangle_{Thermal} = 0 \quad \text{Fluctuations of } \textit{thermal} \text{ nature}$$

The term representing the interference among the modes of the zpf is missing

$$\langle \mathcal{E}^2 \rangle_{Total} = \langle \mathcal{E}_T^2 \rangle + \langle \mathcal{E}_0^2 \rangle + 2 \langle \mathcal{E}_T \mathcal{E}_0 \rangle \rightarrow \langle \mathcal{E}^2 \rangle_{Total} = \langle \mathcal{E}^2 \rangle$$

The distribution W_g cannot account for *all* fluctuations, but only for the *thermal* ones. Thus we look for W_s such that

$$\langle \mathcal{E}^2 \rangle = \langle \mathcal{E}^2 \rangle$$

This generalizes the result:

Relates *thermal* fluctuations with
thermal energy

$$\langle \mathcal{E}^2 \rangle_{Thermal} = \langle \mathcal{E}^2 \rangle_{Equipartition}$$

The distribution that satisfies this requirement and maximizes the entropy is

$$W_s(\mathcal{E}) = \frac{1}{U} e^{-\mathcal{E}/U} \quad \text{It includes fluctuations beyond the thermal ones.}$$



Non-thermal fluctuations

The corresponding entropy is then

$$S_s = -k \int W_s \ln W_s dE = k \ln U$$

Therefore,

$$\frac{S}{U} = \frac{k}{U} \longleftrightarrow \frac{S_{\text{Thermal}}}{U} = \frac{1}{T} \quad \text{Planck and Einstein}$$

Coincide only when the zpe is zero

Thus the existence of the zpe requires a more general (statistical rather than thermodynamic) entropy that should accommodate *all* fluctuations.

Phase space description

The quadratures p, q of the harmonic oscillator are related to its energy according to

$$E = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 q^2$$

Noting that the energy distribution stands for a reduced probability density in the action-angle variables space,

$$W(E) = \int_0^{2\pi} W(\theta, E) d\theta$$



● Non-thermal fluctuations

we can derive the probability density in the phase space (q,p) taking into account that

$$W(E, U) dE dU = \int W(p, q) dp dq$$

The resulting probability density has the form

$$W(p, q) = \frac{\eta_0}{2\pi\hbar} \exp\left[-\frac{p^2 + m^2 \eta_0^2 q^2}{2mU}\right] \quad \text{Wigner function}$$

$$U = E_0 \coth E_0 \quad \text{with } E_0 = \frac{1}{2} \hbar \omega$$

W can be decomposed as the product of two gaussian distributions with variances

$$\sigma_p^2 = mU \quad \sigma_q^2 = U/m \eta_0^2$$

$$\sigma_q^2 \sigma_p^2 = \frac{U^2}{\eta_0^2} = \frac{1}{\eta_0^2} E_0^2 = \frac{\hbar^2}{4} = \frac{E_0^2}{4}$$

$\sigma_p^2 = U^2 \approx E_0^2$ non-thermal fluctuations
 $\sigma_q^2 = \frac{E_0^2}{4}$ thermal fluctuations

Quantum fluctuations are due to the fluctuations of the zpe



● Concluding remarks

- Wien's law allows for a zero-point energy that breaks down with the classical (equipartition) result.
- This zero-point energy **unequivocally** determines Planck's law. No explicit quantum hypothesis is needed.
- The **fluctuating** zero-point energy lies at the root of quantum fluctuations.
- Quantum theory could have been born from this notion (zpe physics).
- Does this physical element is richer?....