

# New light sources based on nonlinear photonic crystal fibers



Miguel V. Andrés

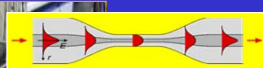



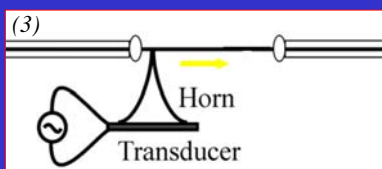
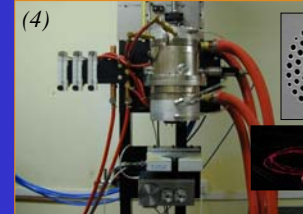
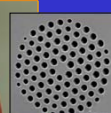

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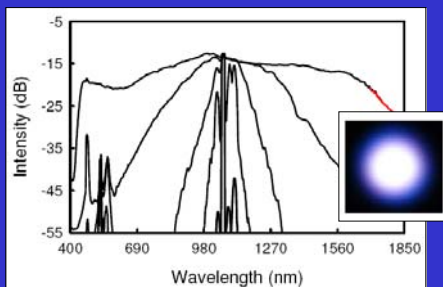


**Fiber-optics components technologies**

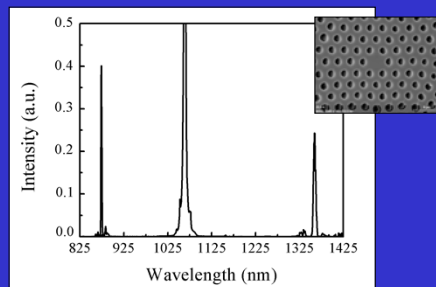
- **Fabrication of fiber-optics components**
  - Optical fiber tapers: fusion & pulling technique
    - (1)   

  - Fiber Bragg gratings and LPG
    - (2)   

  - In-Fiber acousto-optic devices
    - (3) 
  - Microstructured optical fiber
    - (4)   
  

- **Applications:**
  - All-fiber light sources
  - Optical communications and microwave photonics
  - Optical fiber sensors

- The development of optical fiber light sources based on non linear photonic crystal fibers

- Photonic crystal fibers can be designed with specific dispersion properties for an efficient exploitation of in-fiber non linear effects
- Some postprocessing techniques permit an accurate control and a fine adjustment of the dispersion
- Our work has focused on supercontinuum generation and photon pairs generation



Supercontinuum generation in a tapered fiber



Photon pairs generation through FWM

# Supercontinuum and photon pairs generation in photonic crystal fibers

## I. Introduction

- I.1. Fabrication of photonic crystal fibers (PCF)
- I.2. Fibers with special dispersion properties

## II. Supercontinuum generation

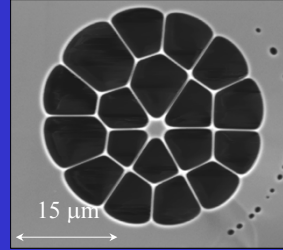
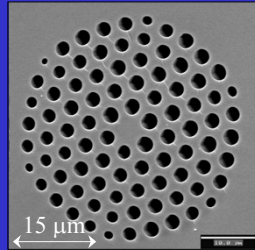
- II.1. Interplay between supercontinuum and dispersion
- II.2. A Y-shaped PCF with a Ge-doped core
- II.3. Er-doped PCF: Shaping the supercontinuum
- II.4. Supercontinuum generation in a tapered PCF

## III. Photon pairs generation through degenerated FWM

- III.1. Pump power induced wavelength shift
- III.2. Fine tuning of FWM bands in Er-doped PCF

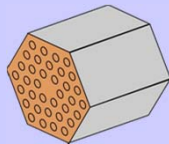
## IV. Conclusions

• Some examples of microstructured optical fibers



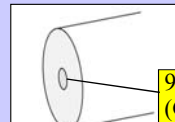
• Guidance mechanism

Microstructured optical fiber



Silica with a microstructure of air holes

Conventional optical fiber



9 μm core (Ge-doped silica)

125 μm cladding (silica)

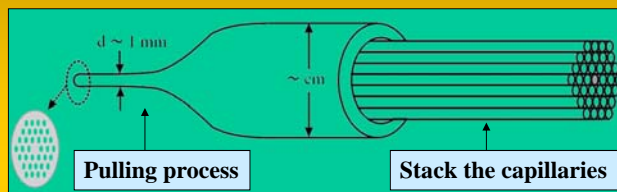
• Fabrication of photonic crystal fibers

First step

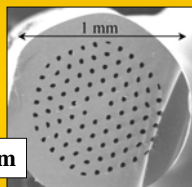


Fabrication of capillaries

Second step

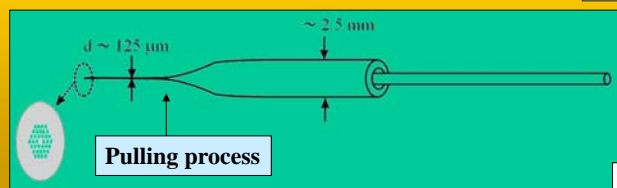


Stack



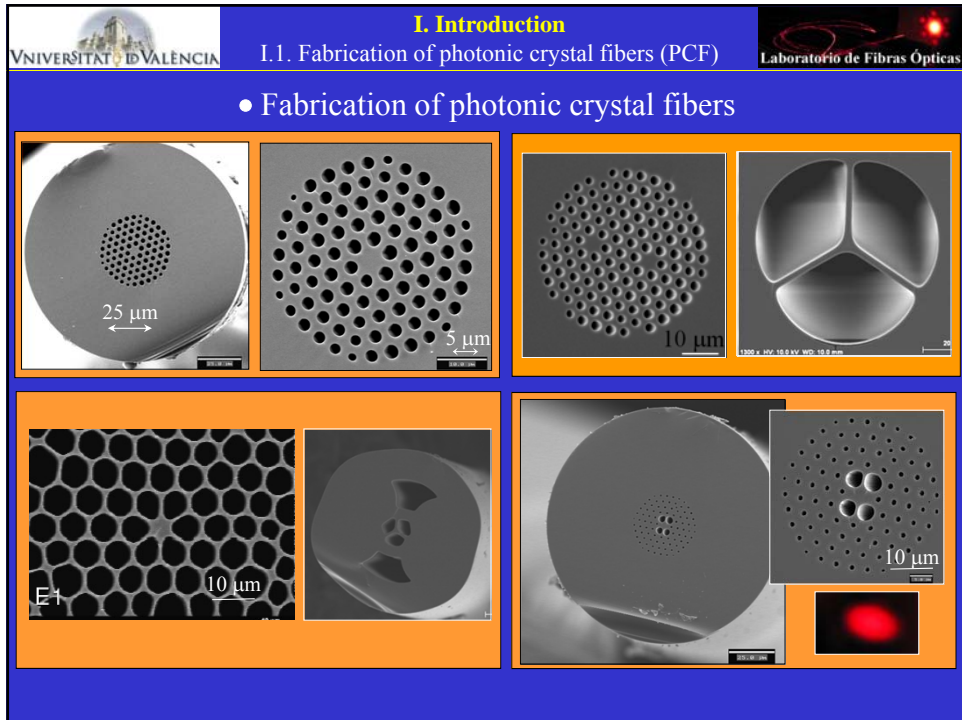
Preform


Third step




Far field







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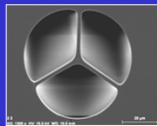
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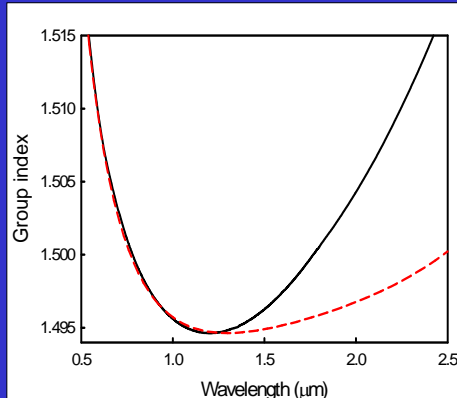
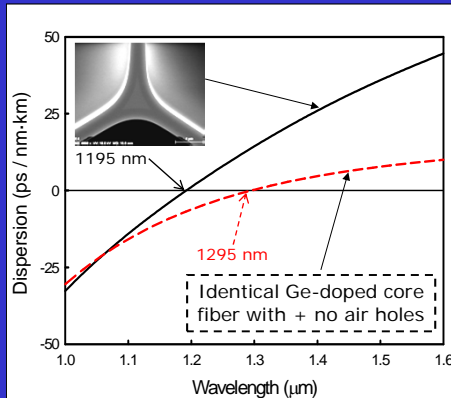
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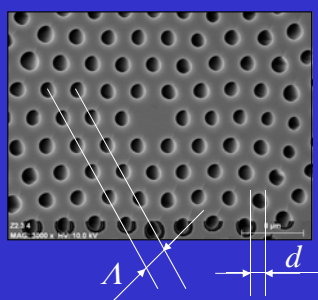
• Silica based photonic crystal fibers with special dispersion properties



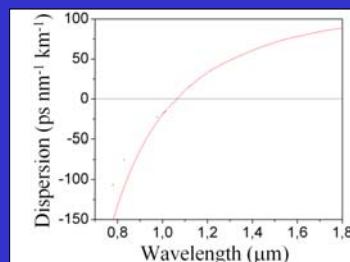
- A highly Ge-doped-core Y-shaped microstructured fiber
- A Ge-doped area of  $6.5 \mu\text{m}^2$  with 20 %w of  $\text{GeO}_2$
- Enhanced non linear characteristics:  
 $n_2 = 1.3 \times n_2(\text{SiO}_2)$ ,  $g_R(12.9 \text{ THz}) = 3 \times g_R(\text{SiO}_2)$



• Silica based photonic crystal fibers with special dispersion properties



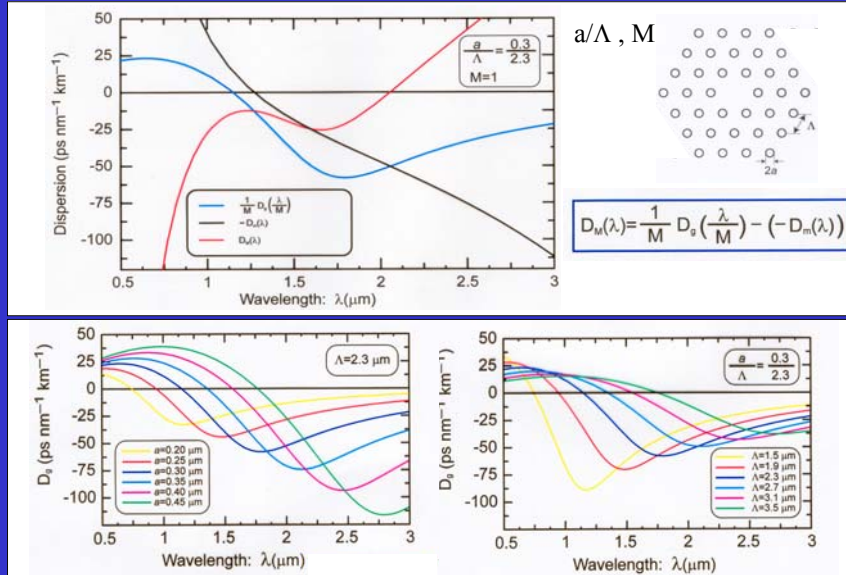
- A set of nonlinear photonic crystal fiber with slightly different zero dispersion wavelength
- Fine adjustment of the dispersion during fabrication



Fiber	Data from SEM images		$\lambda_{ZDW}$ (μm)
	$d$ (μm)	$A$ (μm)	
PCF-A	1.42	3.57	1.082
PCF-B	1.99	4.03	1.070
PCF-C	1.94	3.92	1.065



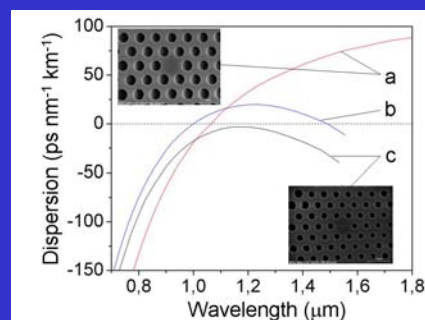
– Designing the dispersion properties of photonic crystal fibers



Opt. Express, pp. 687-697, 2001

- Control of the dispersion properties during fabrication
- A series of PCF produced from a given preform

Fiber	Data from SEM images	
	$d (\mu\text{m})$	$\Lambda (\mu\text{m})$
PCF-a	1.56	2.61
PCF-b	0.72	1.73
PCF-c	0.63	1.70



- PCF-a: One zero dispersion wavelength at  $\sim 1050 \text{ nm}$
- PCF-b: Two zero dispersion wavelengths at  $1000$  and  $1450 \text{ nm}$
- PCF-c: Completely normal dispersion PCF

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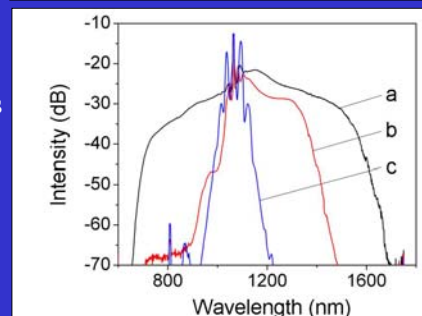
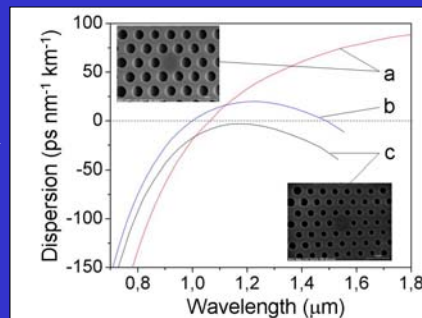
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### IV. Conclusions

### II. Supercontinuum generation

#### II.1. Interplay between supercontinuum and dispersion

- Supercontinuum generation with a pulsed laser: 1064 nm, 0.7 ns
- PCF-a (5 mW)
  - One zero dispersion wavelength at  $\sim 1050$  nm
  - Anomalous dispersion, modulation instability, solitons, self frequency shift (SFS)
  - Dispersive waves
- PCF-b (5 mW)
  - Two zero dispersion wavelengths at 1000 and 1450 nm
  - The second zero dispersion wavelength limits the anomalous dispersion region and prevents the SFS mechanism to extend the spectrum beyond 1450 nm
- PCF-c (7 mW)
  - Completely normal dispersion PCF
  - Raman bands





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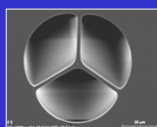
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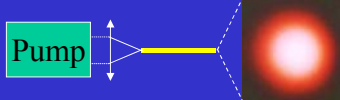
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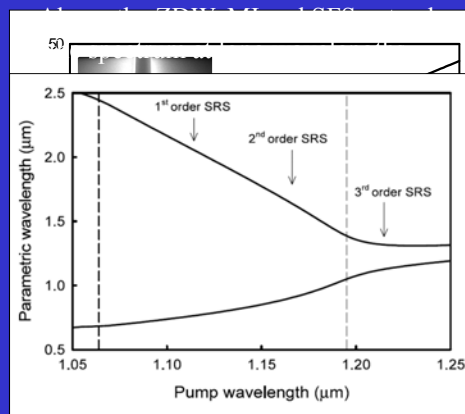
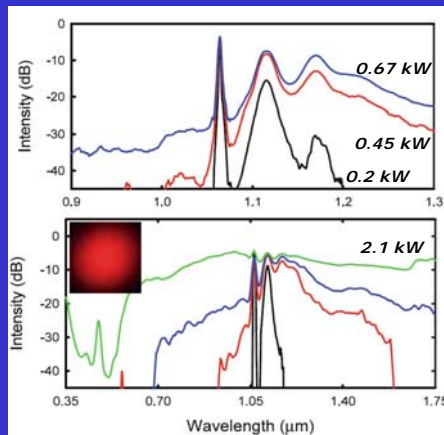
- Supercontinuum generation in a Y-shaped PCF with a Ge-doped core



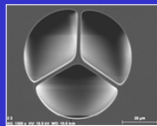
- 1064 nm
- 9 ns pulses
- 5 m length



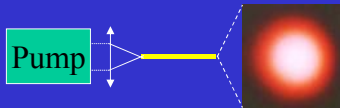
- Large Raman gain



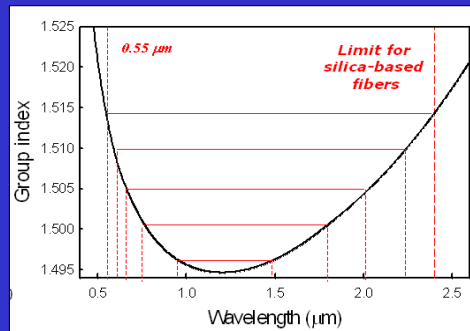
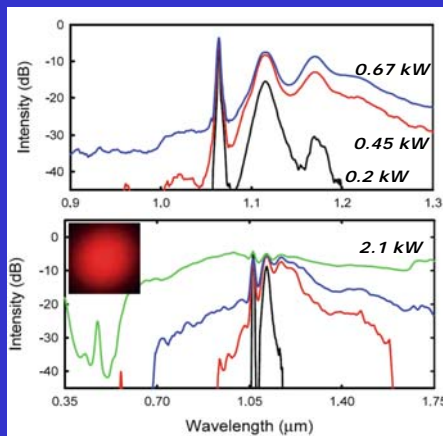
- Supercontinuum generation in a Y-shaped PCF with a Ge-doped core



- 1064 nm
- 9 ns pulses
- 5 m length



- SFS solitons produce dispersive waves that couple power efficiently to shorter wavelengths with the same group index



Appl. Phys. B, pp. 371–376, 2010

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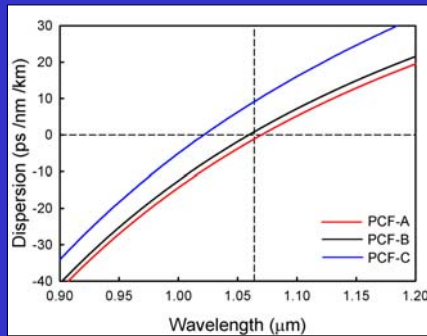
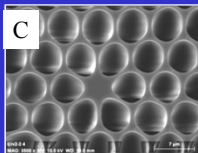
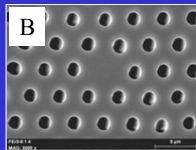
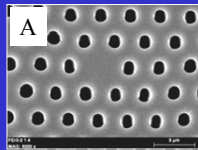
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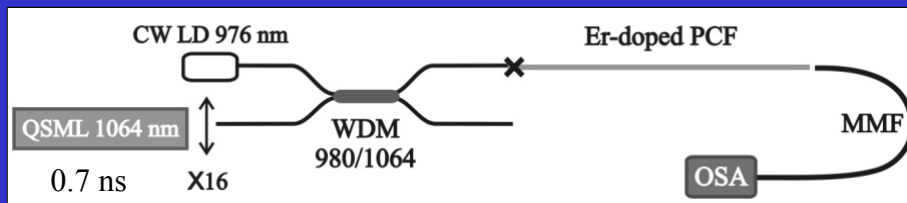
### IV. Conclusions

• Erbium doped PCF

	$A$ ( $\mu\text{m}$ )	$d$ ( $\mu\text{m}$ )	$D_{\text{core}}$ ( $\mu\text{m}$ )	$D_{\text{Br}}$ ( $\mu\text{m}$ )	Er concentration (ions/ $\text{m}^3$ )	ZDW ( $\mu\text{m}$ )
PCF-A	3.7	1.5	5.6	1.6	$1.15 \times 10^{26}$	1.071
PCF-B	3.6	1.5	5.6	1.6	$1.15 \times 10^{26}$	1.058
PCF-C	5.7	0.9	4.5	1.4	$1 \times 10^{25}$	1.024

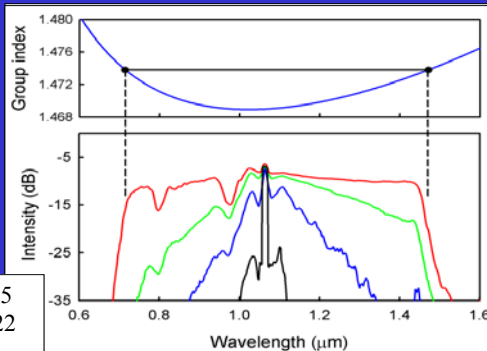


• Erbium doped PCF



- Supercontinuum spectra generated in 3 m of PCF-B (anomalous regime) when  $P_{1064} = 20$  mW
- Black line:  $P_{976} = 0$  mW
- Red line:  $P_{976} = 80$  mW
- Shaping the spectrum with the absorption band of Er at 1530 nm

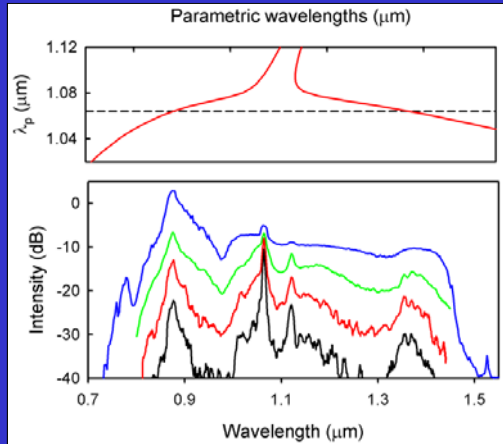
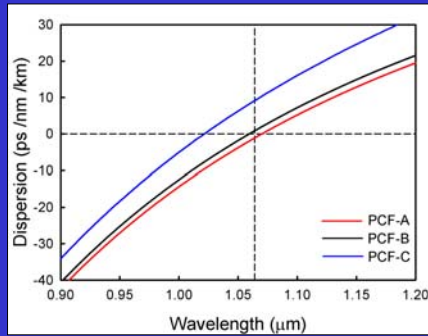
- Supercontinuum spectra generated with 5 m of PCF-B when  $P_{1064} = 4, 10, 10$  and 22 mW



• Erbium doped PCF

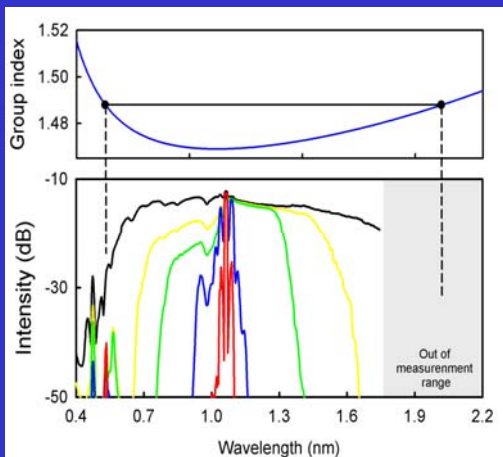
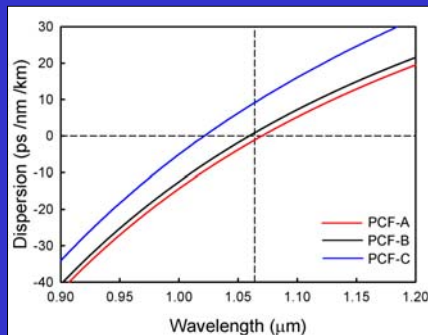


- Supercontinuum spectra generated in 5 m of PCF-A (normal regime) when  $P_{1064} = 4, 10, 22$  and  $30$  mW



• Erbium doped PCF

- Supercontinuum spectra generated in 3 m of PCF-C (anomalous regime) when  $P_{1064} = 8, 16, 30, 44$  and  $66$  mW
- The low concentration of Er does not prevent the SFS mechanism to extend beyond 1530 nm



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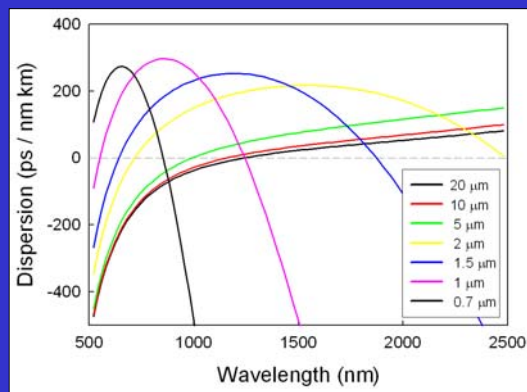
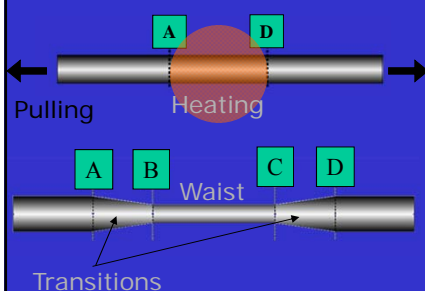
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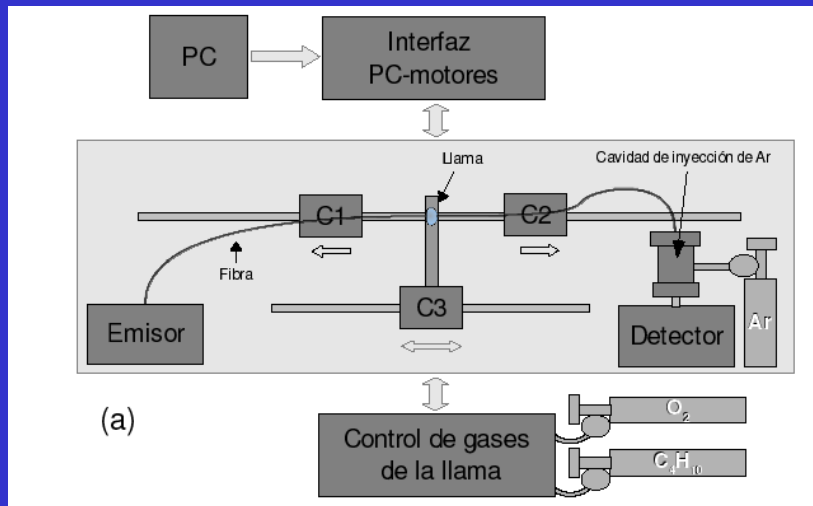
### IV. Conclusions

#### • Tapering PCF



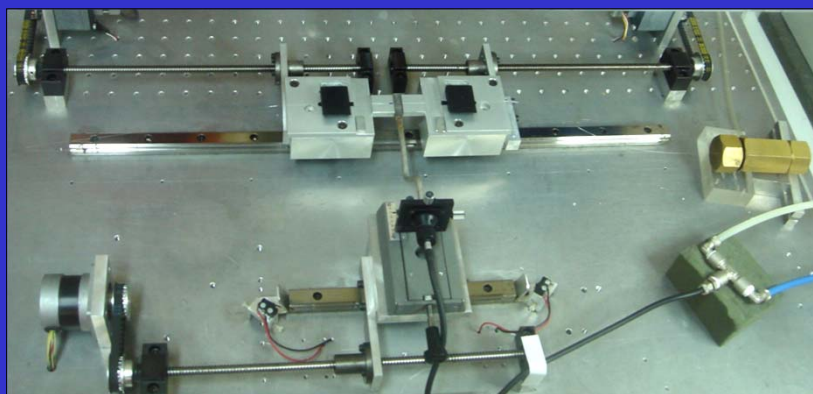
- Tapering a PCF the dispersion can be adjusted with a high degree of freedom
- Along the device the dispersion changes: dispersion engineering
- Small core fibers can be spliced to standard fibers with low losses

• Tapering PCF



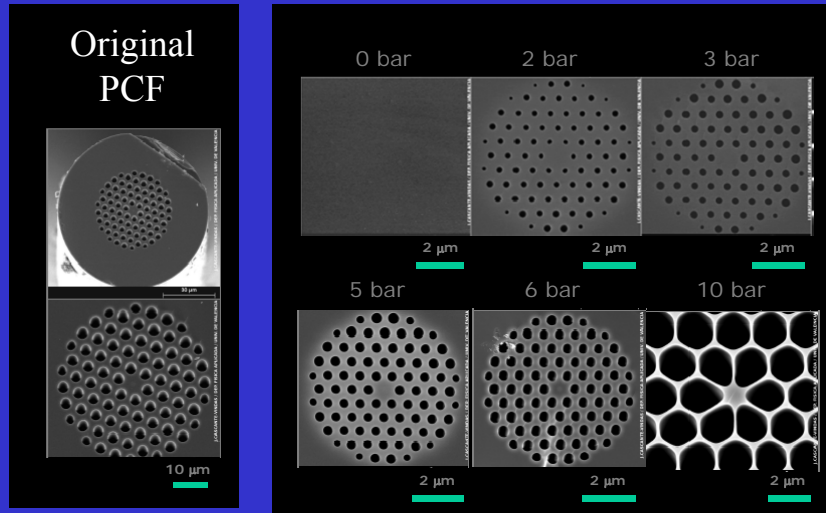
T. A. Birks, Y. W. Li, J. Lightwave Technol. 10, 432-438 (1992)

• Tapering PCF



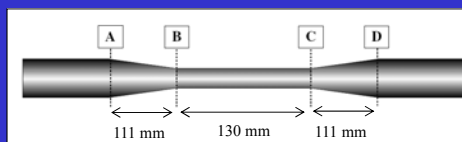
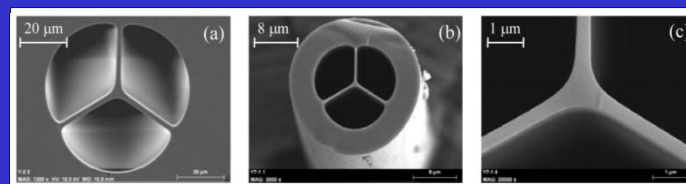


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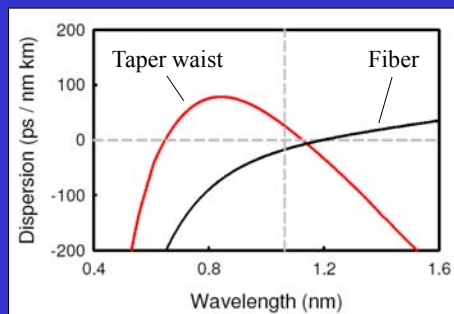


Opt. Commun., pp. 433-438, 2008

• Supercontinuum generation in a tapered Y-shaped PCF



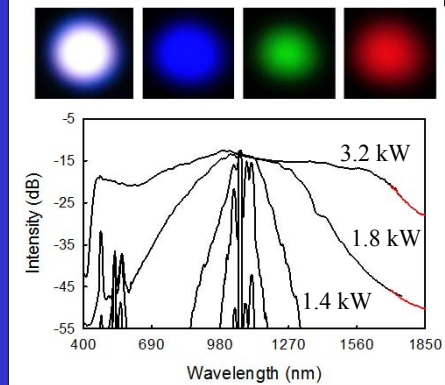
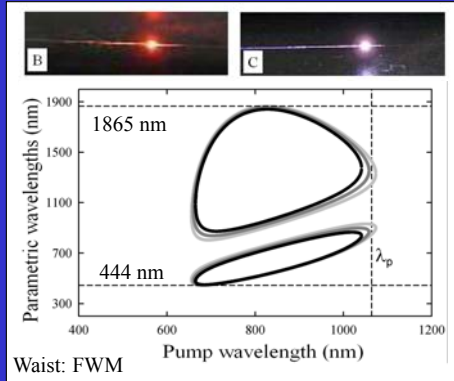
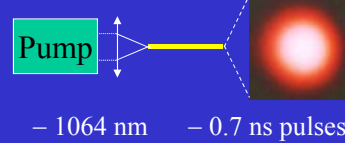
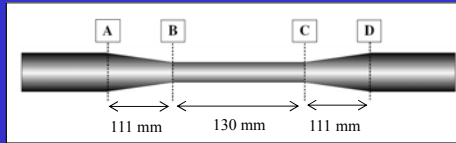
- At the taper waist, the core diameter was  $0.9 \mu\text{m}$
- The first zero dispersion wavelength (ZDW) moves from 1195 to 647 nm along the transitions and, at the taper waist, the fiber exhibits two ZDW at 647 and 1127 nm



## II. Supercontinuum generation

### II.4. Supercontinuum generation in a tapered PCF

#### • Supercontinuum generation in a tapered Y-shaped PCF

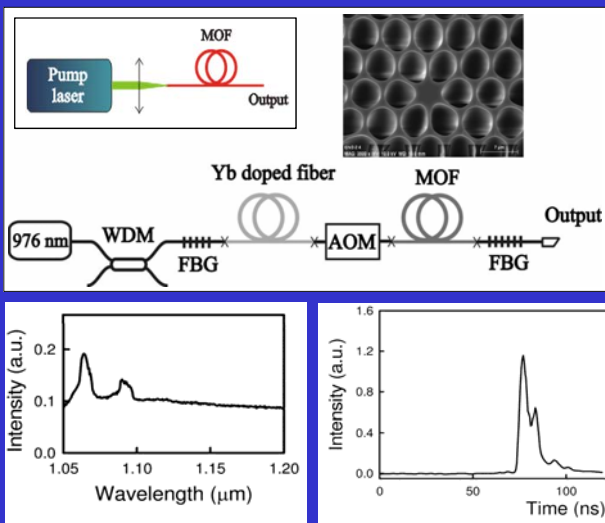
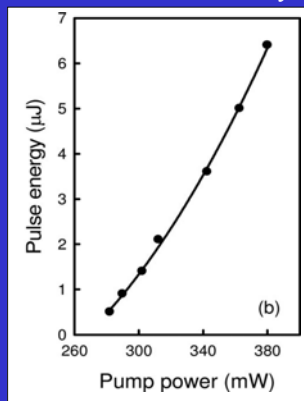


Opt. Express, pp. 14535-14540, 2010

## Supercontinuum generation

#### • Supercontinuum generation using nanosecond pulses

- All-fiber systems



Opt. Lett., pp. 3628-3630, 2009

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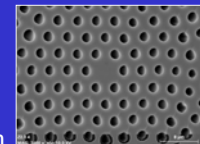
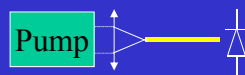
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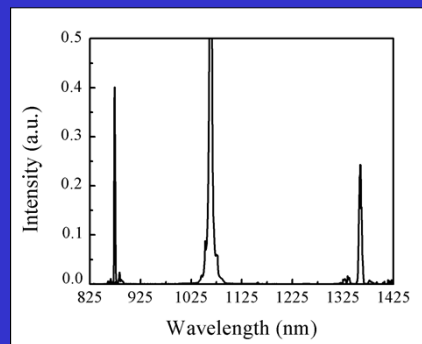
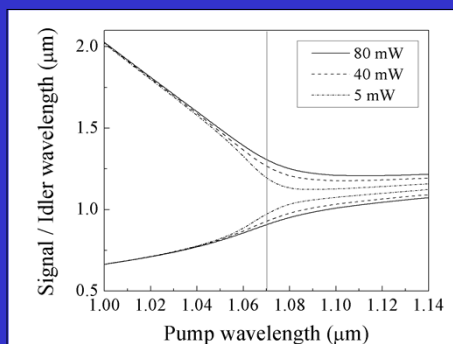
- 1064 nm
- 0.72 ns pulses
- Power average < 160 mW
- 20 kHz repetition rate

$$2 \cdot \omega_p = \omega_i + \omega_s$$

$$\kappa = 2 \cdot \beta(\omega_p) - \beta(\omega_i) - \beta(\omega_s) - 2 \cdot \gamma \cdot P$$

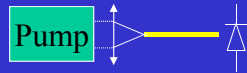


- Fiber length: 0.7 - 1 m
- Fibers with slightly different ZDW



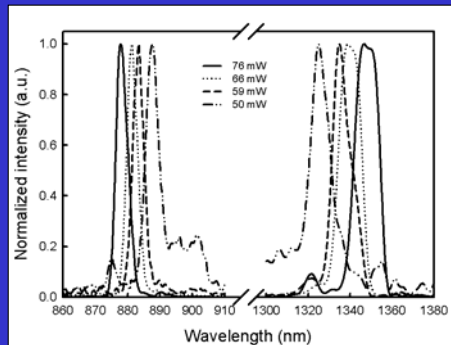
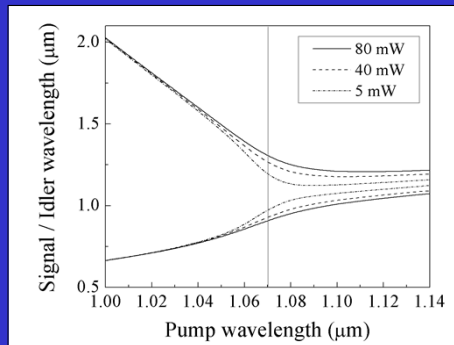
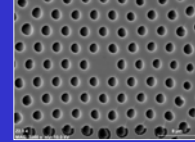
• Photon pairs generation through degenerated FWM

- 1064 nm
- 0.72 ns pulses
- Power average < 160 mW
- 20 kHz repetition rate

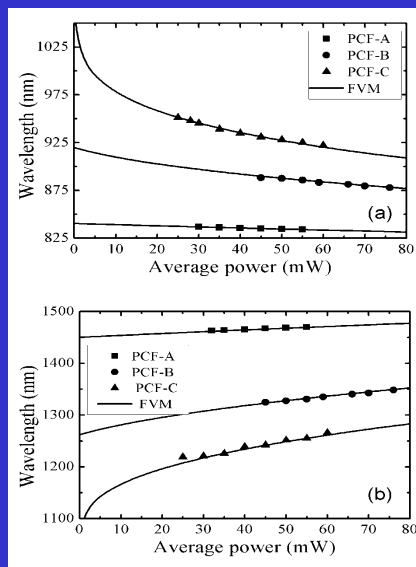


– Fiber length: 0.7 - 1 m

– Fibers with slightly different ZDW



• Photon pairs generation through degenerated FWM



Fiber	SEM images		Best fit		$\lambda_{ZDW}$ (μm)
	$d$ (μm)	$\Lambda$ (μm)	$d$ (μm)	$\Lambda$ (μm)	
PCF-A	1.42	3.57	1.410	3.531	1.082
PCF-B	1.99	4.03	1.864	3.771	1.070
PCF-C	1.94	3.92	1.771	3.589	1.065

- The linear properties of the fibers were modeled using a fully vectorial method.
- For PCF-C, the signal and idler bands shifted from 951 nm to 922 nm, and from 1219 nm to 1266 nm, respectively, in a pump power range 35 mW.
- For a 80 mW pump power range, the theoretical calculations predict signal and idler bands shifts of 87 nm and 130 nm.

Photon. Technol. Lett., pp. 1010-1012, 2011

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- II.4. Supercontinuum generation in a tapered PCF

### III. Photon pairs generation through degenerated FWM

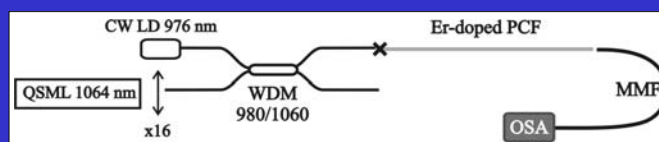
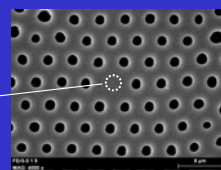
- III.1. Pump power induced wavelength shift
- III.2. Fine tuning of FWM bands in Er-doped PCF

### IV. Conclusions

### III. Photon pairs generation through FWM

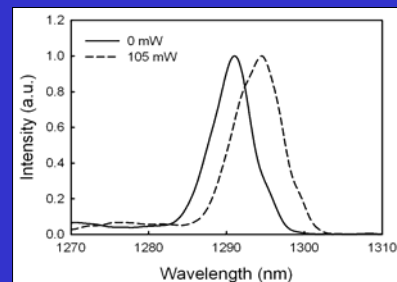
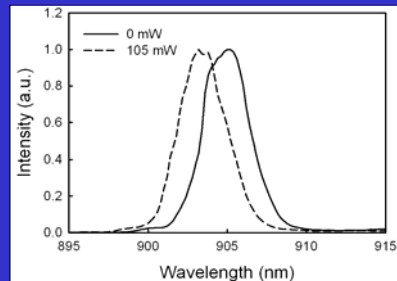
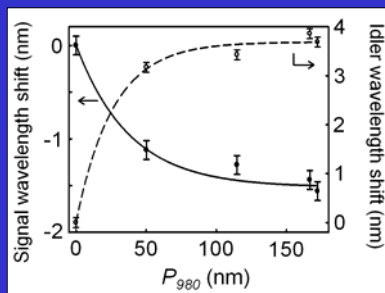
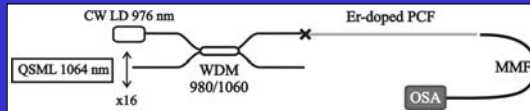
#### III.2. Fine tuning of FWM bands in Er-doped PCF

- Normal dispersion Er-doped fiber at 1064 nm: dynamic control by pumping the fiber with a CW 976 nm laser diode
  - Pitch  $\Lambda = 3.72 \mu\text{m}$
  - Hole diameter  $d = 1.46 \mu\text{m}$
  - Core diameter  $d_c = 5.55 \mu\text{m}$
  - Doped region:  $2.23 \mu\text{m}$
  - Er concentration: 2300 ppm
  - Zero dispersion wavelength  $\lambda_{ZDW} = 1072 \text{ nm}$



- Q-switched Microchip Nd-YAG laser: 1064 nm, 0.7 ns, 20 kHz, 160 mW
- Pump power coupled to the PCF: 29 mW
- Pigtailed laser diode: 976 nm
- Maximum power coupled to the PCF: 175 mW

- The CW pump at 976 nm has several effects on the PCF that might contribute to the shift of FWM bands:
  - Change of the nonlinear coefficient  $n_2$
  - Change of refractive index produced by absorption changes (Kramers-Kronig)
  - Thermal effects



- Theoretical discussion: non linear refractive index  $n_2$

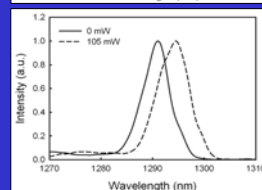
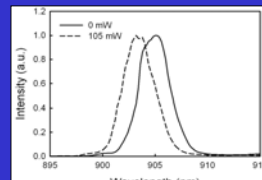
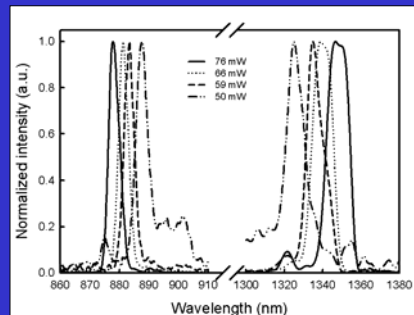
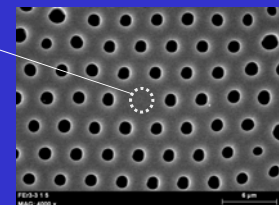
$$2 \cdot \beta(\omega_p) = \beta(\omega_i) + \beta(\omega_s) + 2 \cdot \gamma \cdot P$$

- A small change of the nonlinear refractive index  $n_2$  could be, in principle, responsible for the reported shifts of the FWM bands, but  $n_2$  decreases when pumping the Er-doped fiber\* which would shift the bands towards the pump wavelength

\* H. García et al., Opt. Lett., pp. 1261-1263, 2005

- A decrease of  $n_2$  is equivalent to a decrease of  $P$ :

Er-doped region





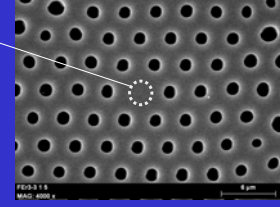
- Theoretical discussion: refractive index changes

– The thermal effects

– First order of perturbation calculus gives:

$$\left( \frac{\partial n_{eff}}{\partial n} \right) \cong 2 \frac{A_{Er}}{A_{eff}}$$

Er-doped region



$$\left. \begin{aligned} \delta n_s &\cong \delta n_i \cong \delta n_p \\ A_{eff} &\cong \text{constant with } \lambda \\ 2\omega_p &= \omega_s + \omega_i \end{aligned} \right\} \Rightarrow \delta \lambda_i = 0$$

$$\frac{2\pi c}{\lambda_i^2} (n_{gs} - n_{gi}) \delta \lambda_i = \omega_s \left( \frac{\partial n_{eff}}{\partial n} \right)_{\omega_s} \delta n_s + \omega_i \left( \frac{\partial n_{eff}}{\partial n} \right)_{\omega_i} \delta n_i - 2\omega_p \left( \frac{\partial n_{eff}}{\partial n} \right)_{\omega_p} \delta n_p$$

– Using an accurate calculation of the derivatives:

$\Delta n_g$	$(\partial n_{eff} / \partial n_{er})_{\omega_i}$	$(\partial n_{eff} / \partial n_{er})_{\omega_p}$	$(\partial n_{eff} / \partial n_{er})_{\omega_s}$
$3.8 \times 10^{-4}$	0.45	0.49	0.52

and taking into account the experimental value  $\delta \lambda_i = 3.7 \text{ nm}$  ( $P_{976} = 160 \text{ mW}$ ) we can estimate the refractive index increment that would be required:  $\delta n = 4.7 \times 10^{-6}$

This increment of the refractive index in the Er-doped region, with respect to the surrounding glass, corresponds to a temperature difference of about 0.55 K, which would require a pump power 25 times higher ( $\sim 4 \text{ W}$ )

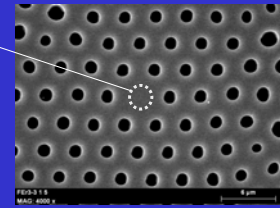
- Theoretical discussion: refractive index changes

– Which refractive changes would enhance the shift of the parametric bands?

– Taking into account that:

$$\left( \frac{\partial n_{eff}}{\partial n} \right) \cong 2 \frac{A_{Er}}{A_{eff}}$$

Er-doped region



$$\frac{2\pi c}{\lambda_i^2} (n_{gs} - n_{gi}) \delta \lambda_i \approx \left( \frac{\partial n_{eff}}{\partial n} \right)_{\omega_s} (\omega_s \delta n_s + \omega_i \delta n_i - 2\omega_p \delta n_p)$$

Therefore, an interesting case is  $\delta n_s > 0$ ,  $\delta n_i > 0$  and  $\delta n_p < 0$

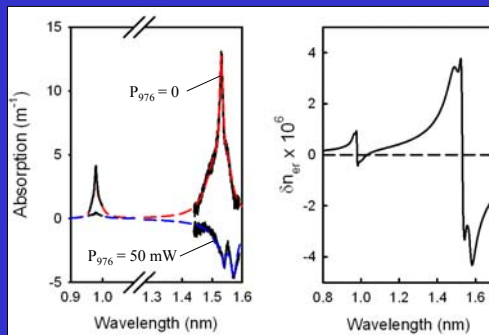
In fact, this is close to the situation that corresponds to the resonant refractive index changes produced by absorption saturation of the 980 nm band of the Erbium ions.

– Using Kramers-Kronig relations:

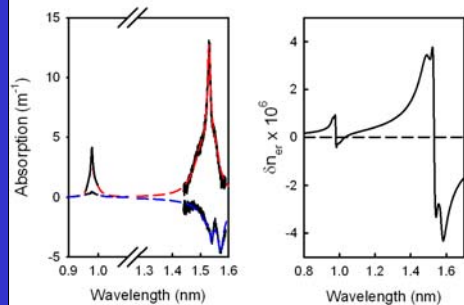
$$\delta n_s = 0.6 \times 10^{-6} \text{ at } 880 \text{ nm}$$

$$\delta n_i = 1.2 \times 10^{-6} \text{ at } 1367 \text{ nm}$$

$$\delta n_p = 1.0 \times 10^{-7} \text{ at } 1064 \text{ nm}$$



- Theoretical discussion: resonant refractive index changes



- Resonant change of refractive index produced by absorption changes (Kramers-Kronig):

$$\delta n_s = 0.6 \times 10^{-6} \text{ at } 880 \text{ nm}$$

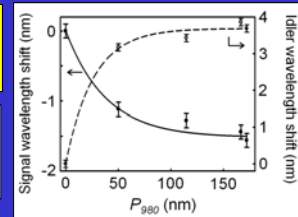
$$\delta n_i = 1.2 \times 10^{-6} \text{ at } 1367 \text{ nm}$$

$$\delta n_p = 1.0 \times 10^{-7} \text{ at } 1064 \text{ nm}$$

- Estimated shifts  $\delta \lambda_s = 1.3 \text{ nm}$  and  $\delta \lambda_i = 3.2 \text{ nm}$  are in good agreement with the experimental values

$$\Delta n_g \delta \omega_i = \omega_s \left( \frac{\partial n_{\text{eff}}}{\partial n} \right)_{\omega_i} \delta n_s + \omega_i \left( \frac{\partial n_{\text{eff}}}{\partial n} \right)_{\omega_i} \delta n_i - 2\omega_p \left( \frac{\partial n_{\text{eff}}}{\partial n} \right)_{\omega_p} \delta n_p$$

$\Delta n_g$	$(\partial n_{\text{eff}} / \partial n_{\text{er}})_{\omega_i}$	$(\partial n_{\text{eff}} / \partial n_{\text{er}})_{\omega_p}$	$(\partial n_{\text{eff}} / \partial n_{\text{er}})_{\omega_s}$
$3.8 \times 10^{-4}$	0.45	0.49	0.52



Opt. Lett., pp. 1226-1228, 2012

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- III.1. Pump power induced wavelength shift
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### ➡ IV. Conclusions

- The photonic crystal fiber technology is crucial for the development of novel fiber light sources based on non linear effects
- Fibers with special dispersion properties can be obtained with this technology and, in addition, the tapering technique permits an accurate engineering of the dispersion properties along short sections of the fiber
- Applications: supercontinuum and photon pairs generation

